

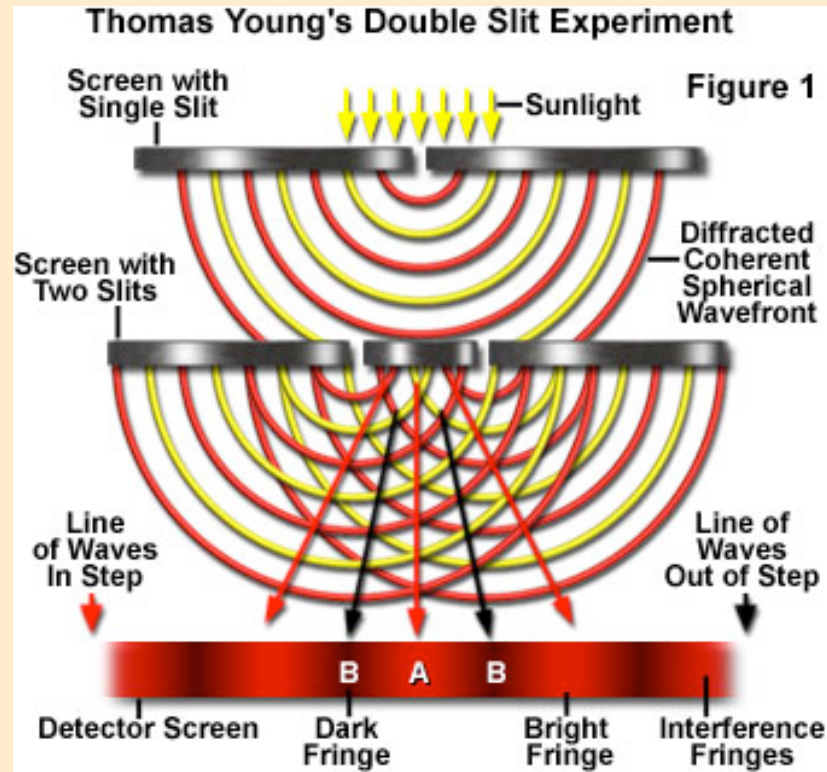
Neutron Interferometry

F. E. Wietfeldt
Tulane University

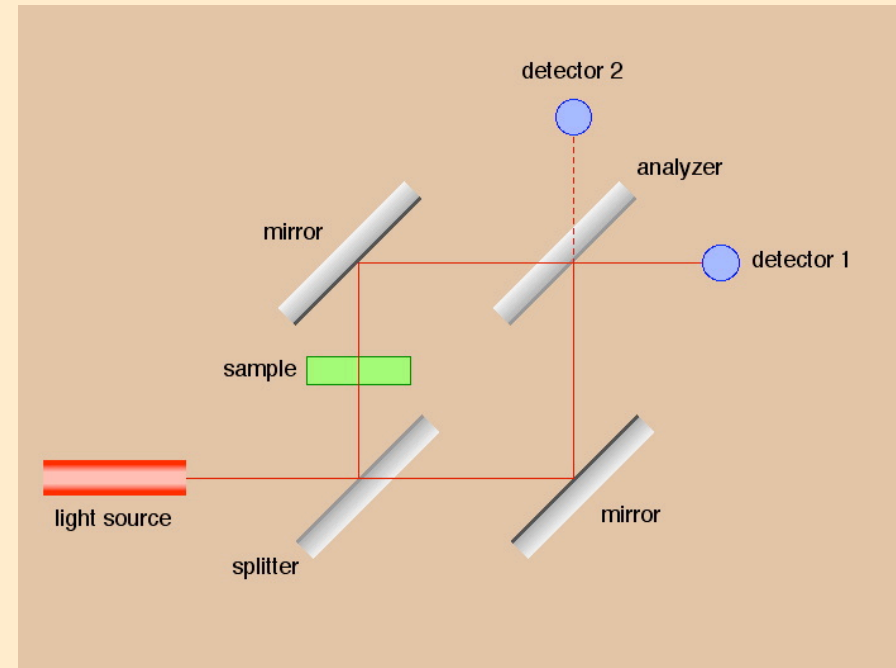
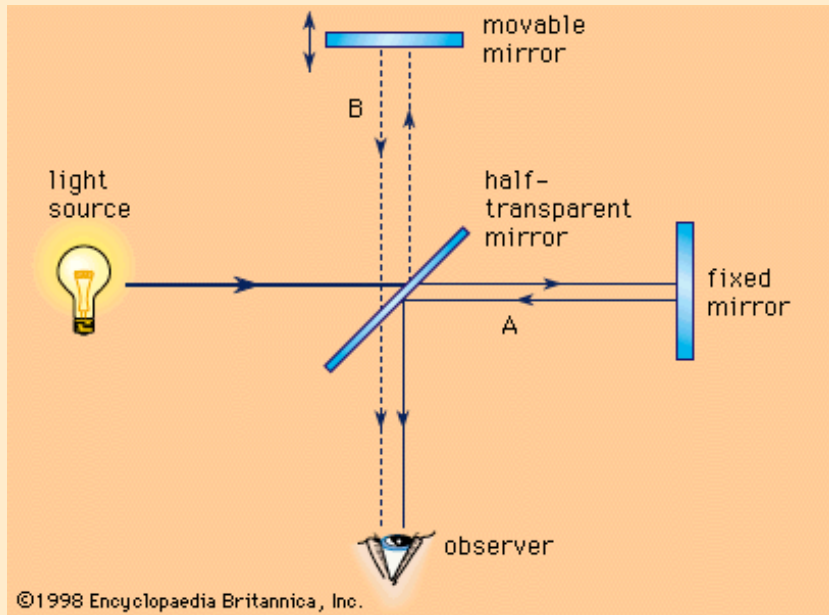
First Summer School on Fundamental Neutron Physics
June 4 – 10, 2006

What is an interferometer?

An interferometer is a device that measures the relative phase between coherent waves by detecting the interference pattern.

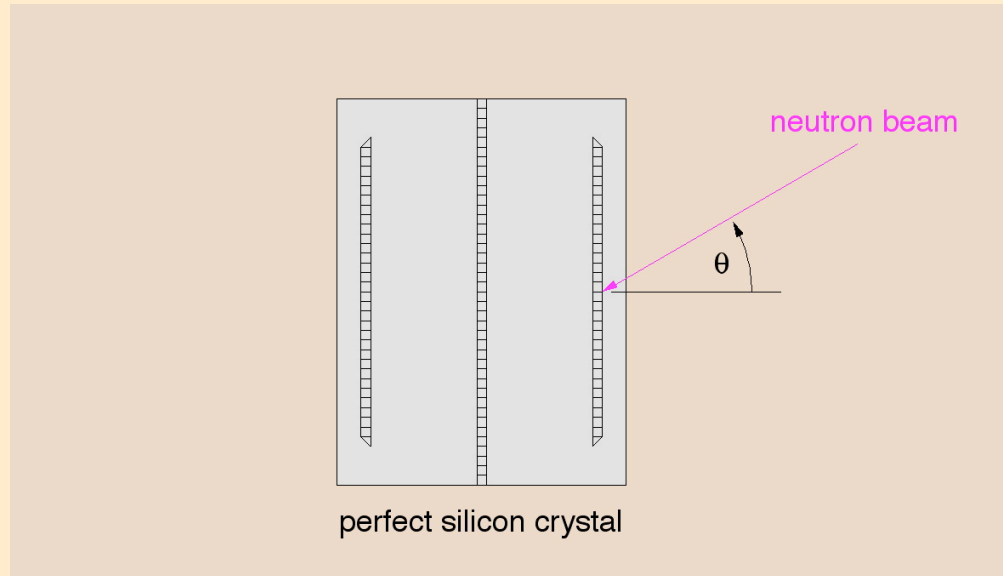


Michelson Interferometer



Mach-Zender Interferometer

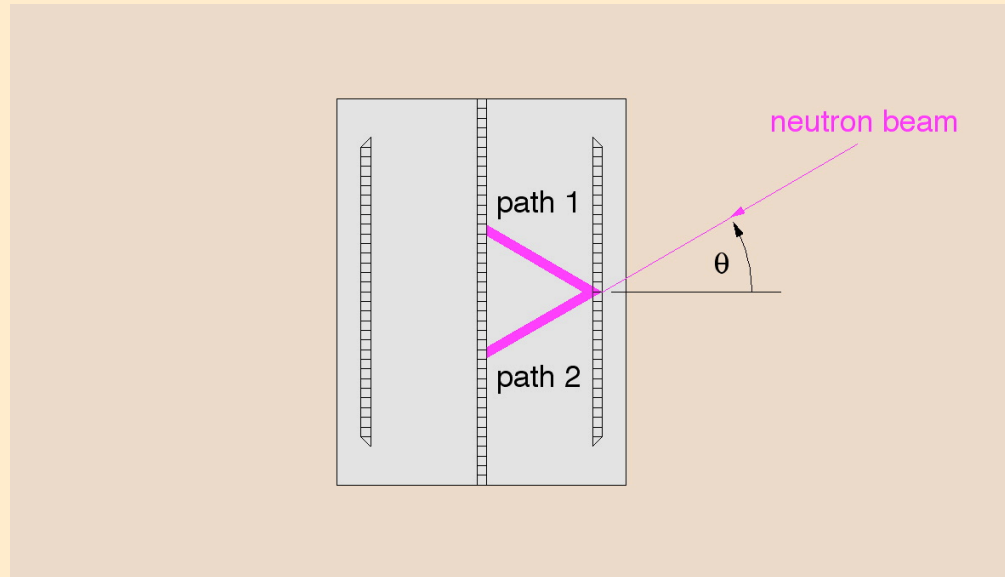
Perfect Crystal LLL Neutron Interferometer



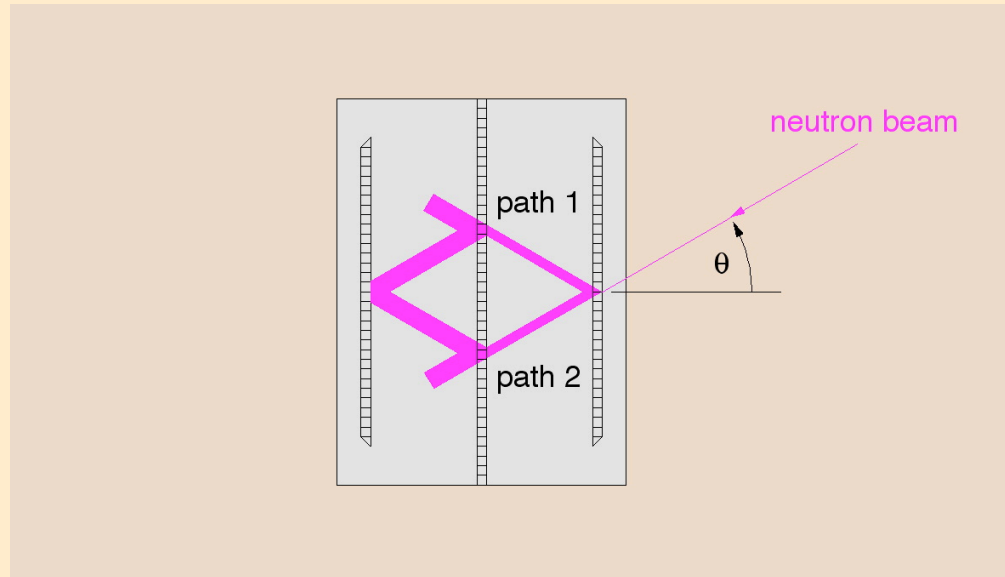
Bragg condition: $n\lambda = 2d \sin \theta$

d = lattice spacing

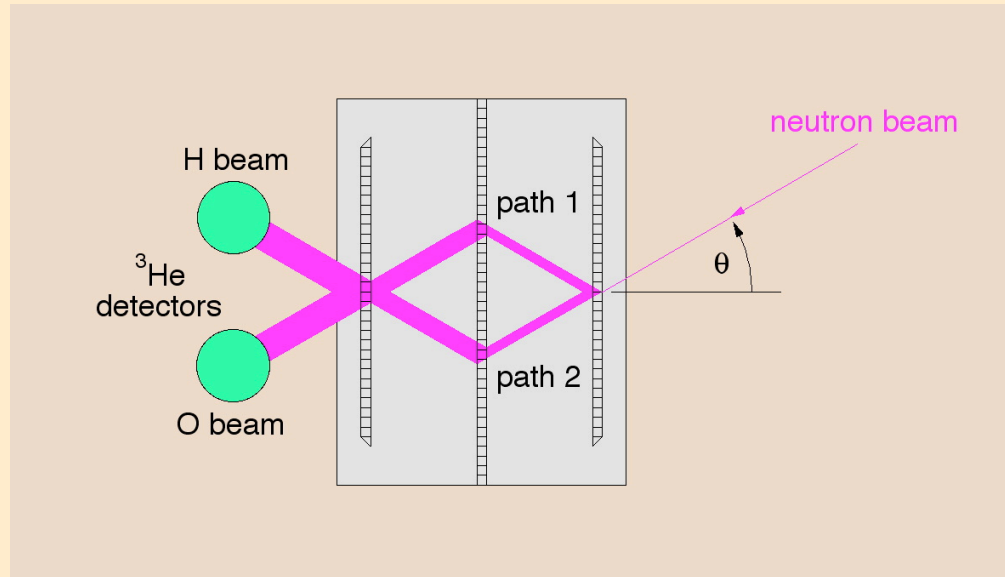
Perfect Crystal LLL Neutron Interferometer



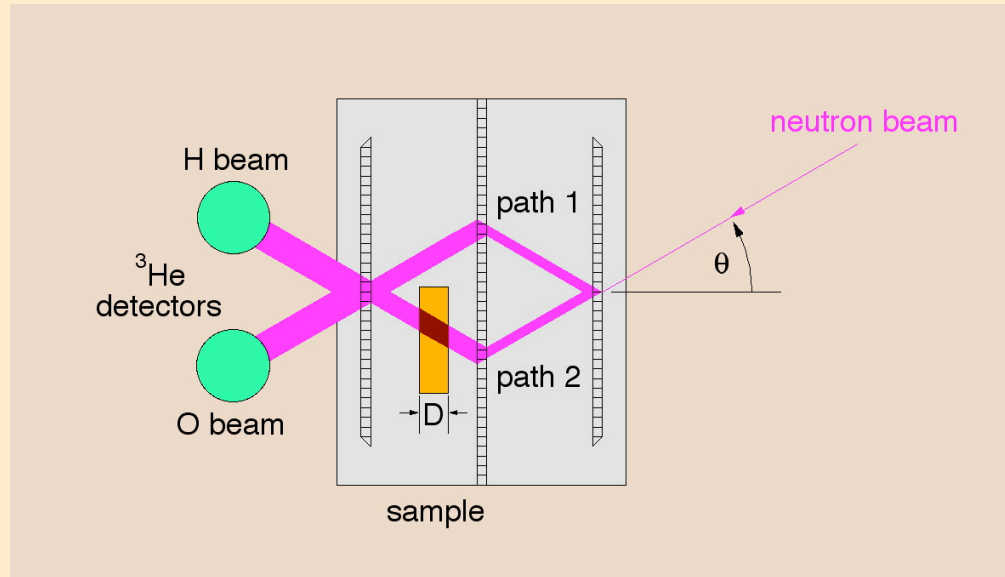
Perfect Crystal LLL Neutron Interferometer



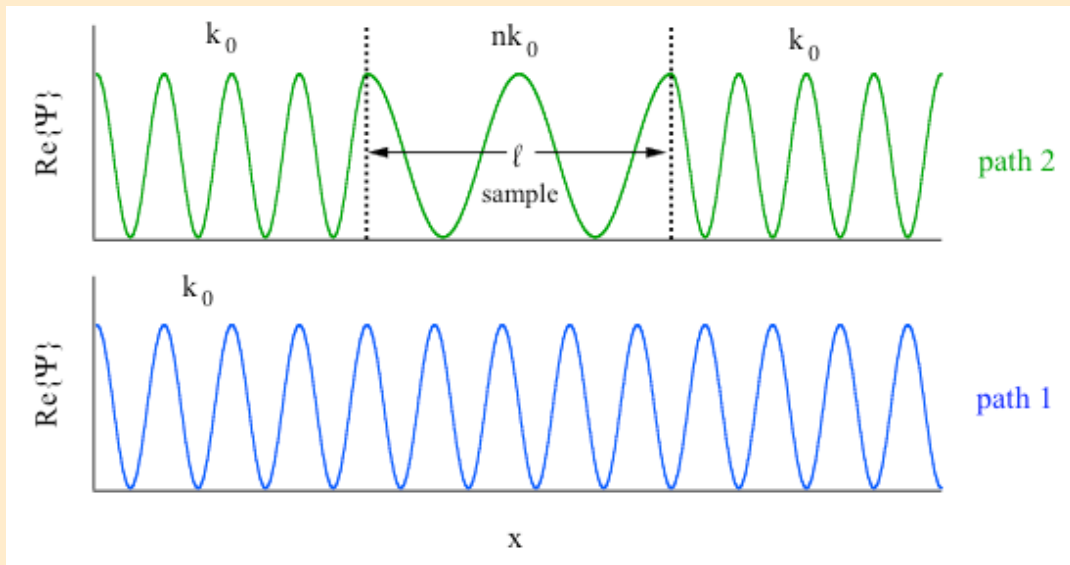
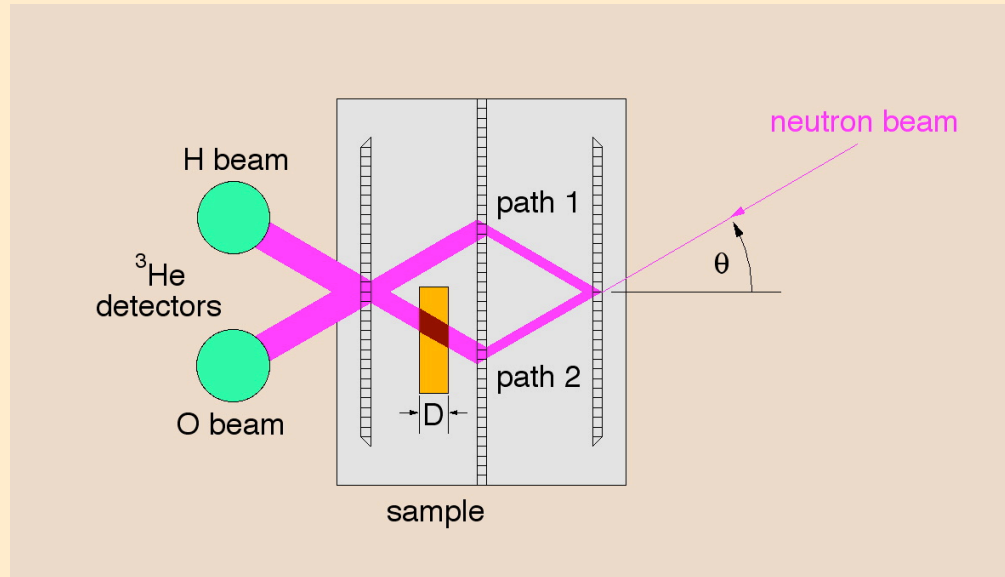
Perfect Crystal LLL Neutron Interferometer



Nuclear Phase Shift



Nuclear Phase Shift

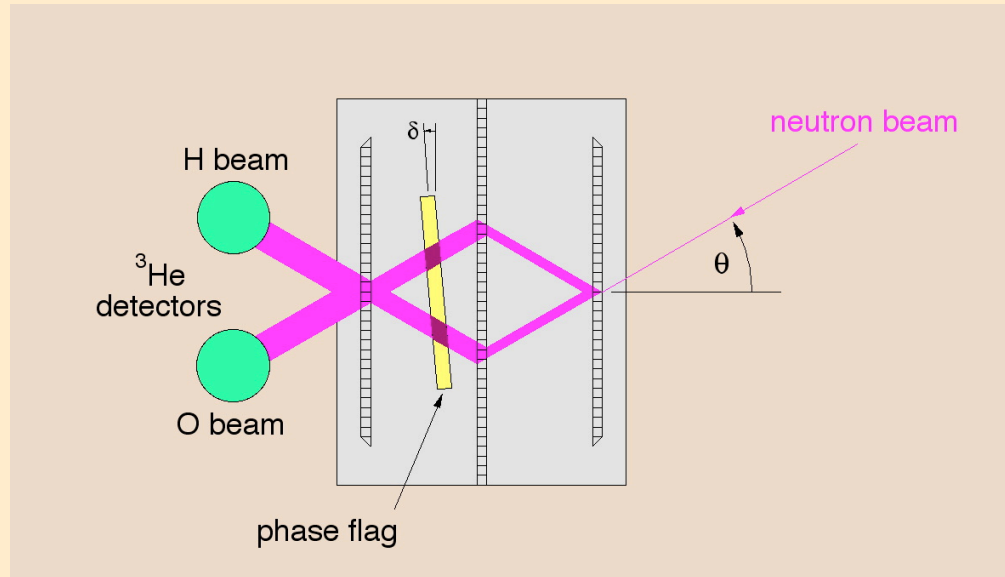


index of refraction: $n = 1 - \frac{Nb\lambda^2}{2\pi}$

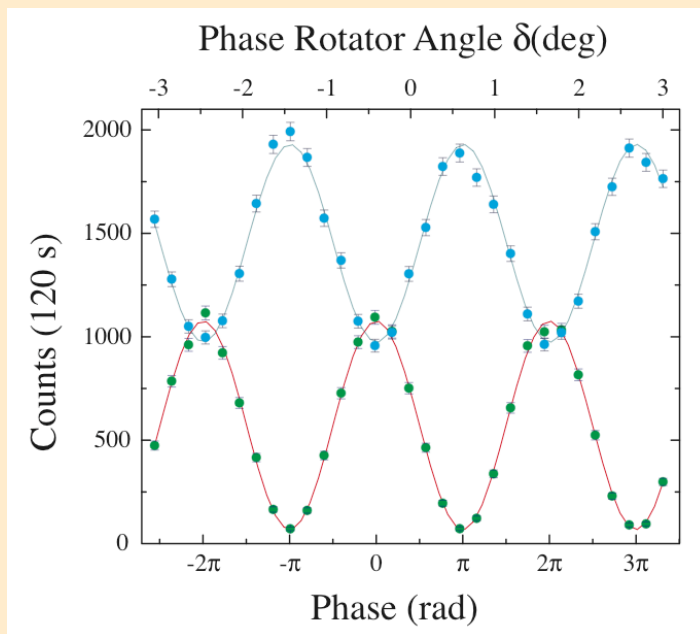
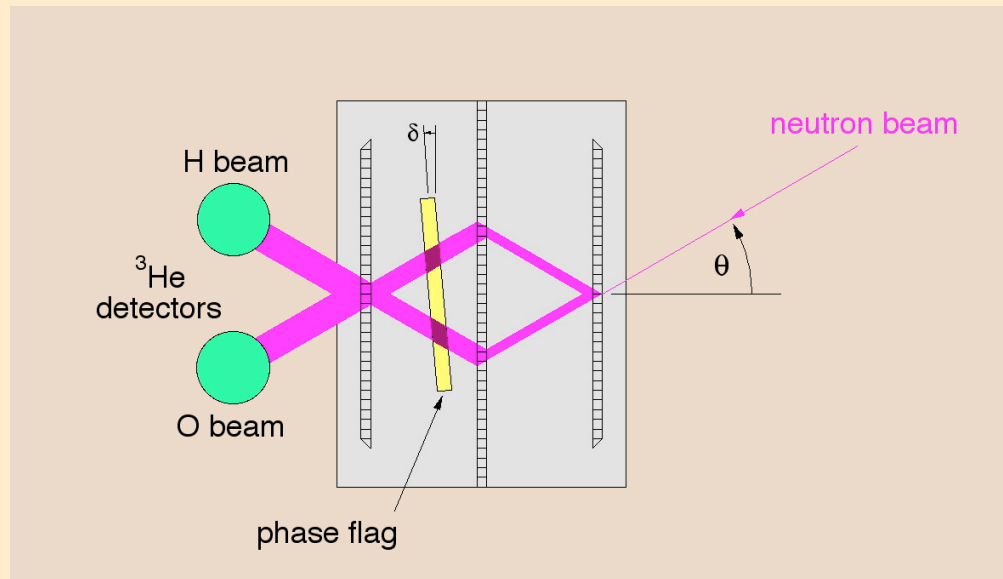
relative phase shift:

$$\Delta\chi = k_0\ell - nk_0\ell = Nb\lambda \frac{D}{\cos\theta}$$

Interferogram



Interferogram

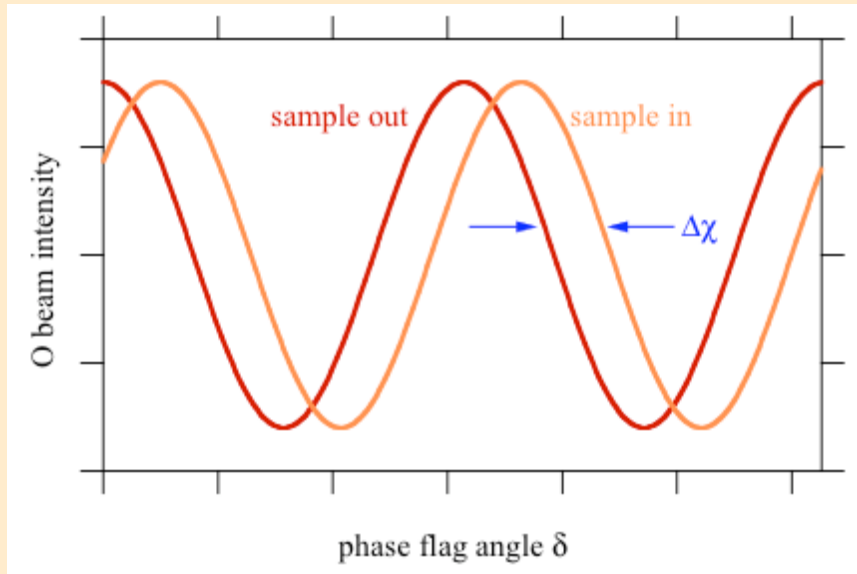
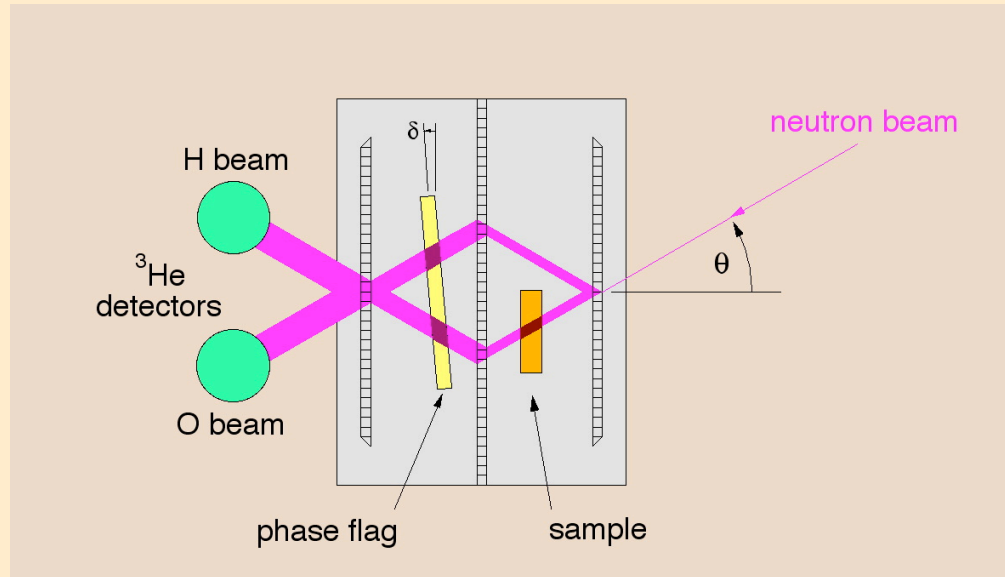


$$\text{O beam: } I_O = A[1 + f \cos(\chi_2 - \chi_1)]$$

$$\text{H beam: } I_H = B - Af \cos(\chi_2 - \chi_1)$$

$$\text{contrast } f = \frac{C_{\max} - C_{\min}}{C_{\max} + C_{\min}} \quad (\text{O-beam})$$

Precision Phase Shift Measurement



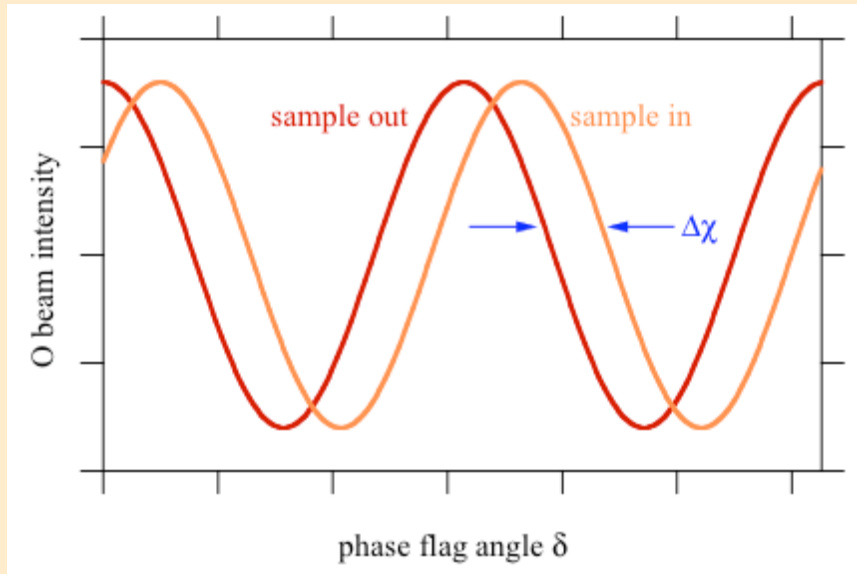
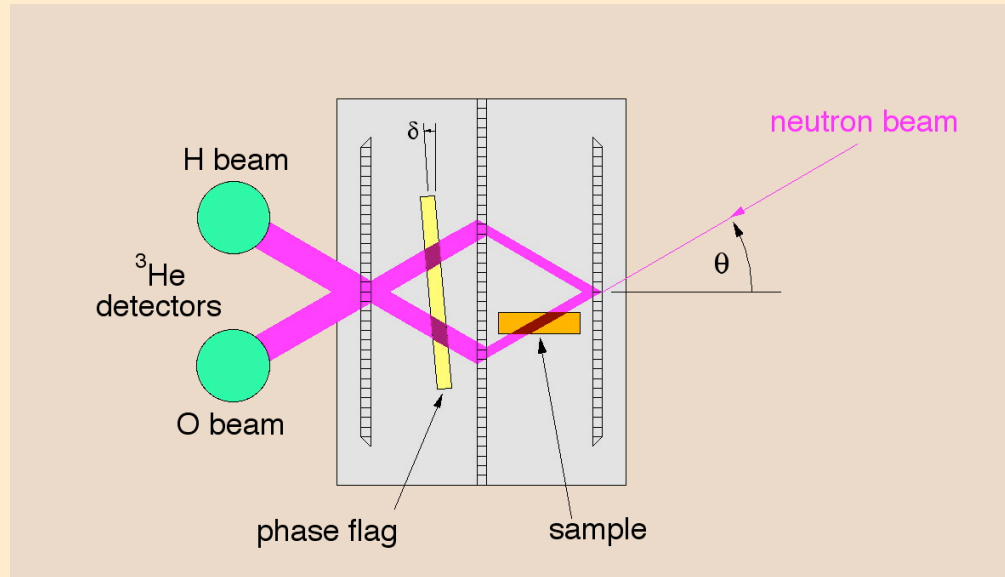
$$\Delta\chi = Nb\lambda \frac{D}{\cos\theta}$$

Example:

aluminum sample, $\lambda = 2.70 \text{ \AA}$, $\langle 111 \rangle$ reflection

$$D = 100 \text{ \mu m} \Rightarrow \Delta\chi = 2\pi$$

Non-Dispersive Geometry

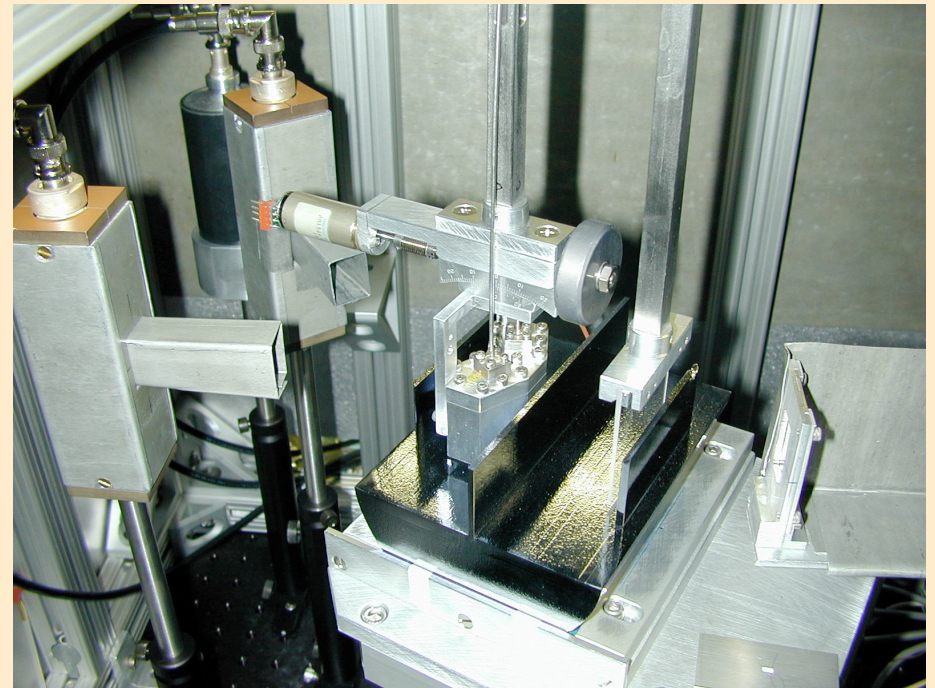
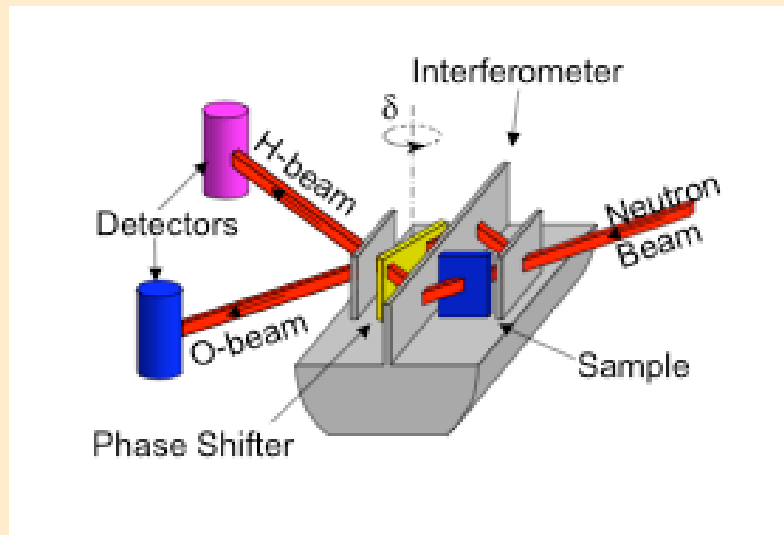


$$\text{path length } \ell = \frac{D}{\sin \theta}$$

$$\Delta\chi = 2Nb d D$$

independent of λ

Perfect Crystal LLL Neutron Interferometer



Precision neutron-interferometric measurement of the coherent neutron-scattering length in silicon

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²St. Petersburg Nuclear Physics Institute, Gatchina, Leningrad District 188350, Russia

³National Institute of Standards and Technology, Gaithersburg, Maryland 20899

⁴Nuclear Physics Institute of CAS, 20568 Rez, Czech Republic

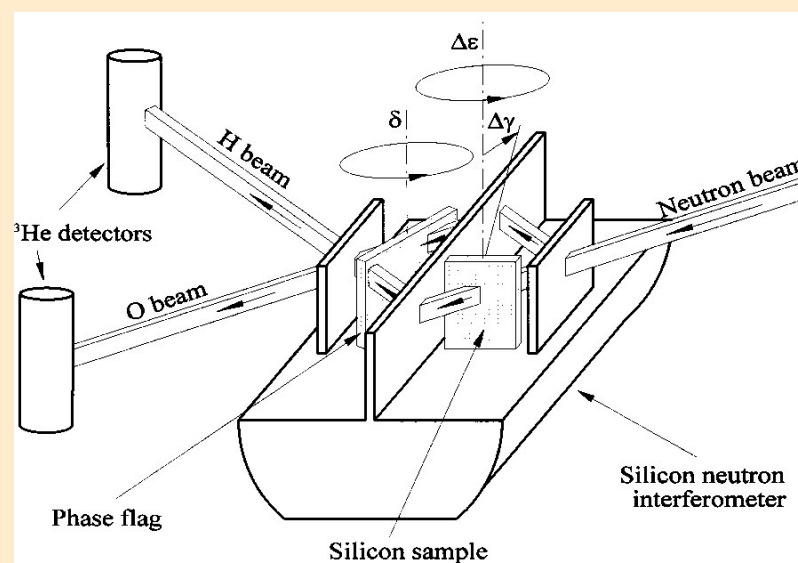
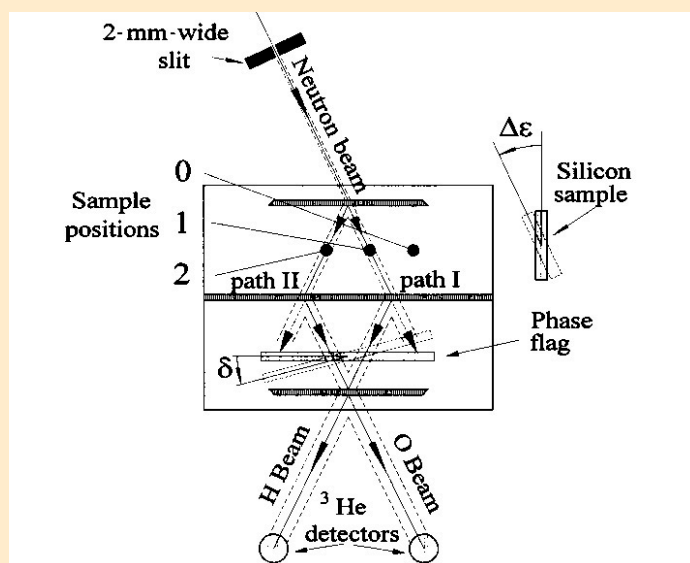
⁵Department of Physics and Astronomy, University of Missouri-Columbia, Columbia, Missouri 65211

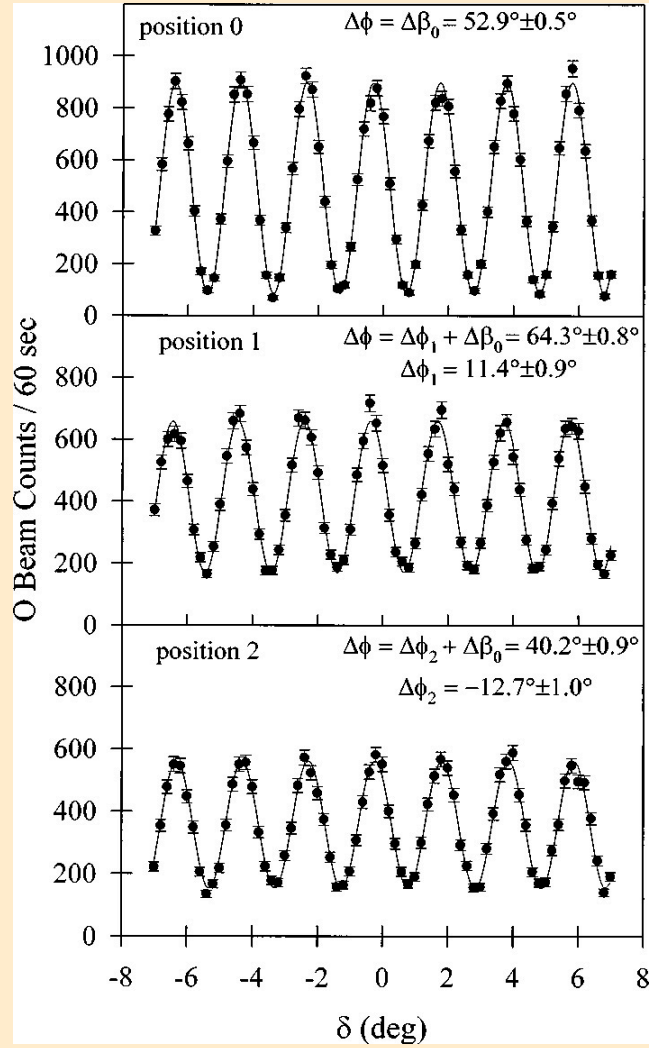
⁶Los Alamos National Laboratory, Los Alamos, New Mexico 87545

(Received 15 August 1997)

The neutron-interferometry (NI) technique provides a precise and direct way to measure the bound, coherent scattering lengths b of low-energy neutrons in solids, liquids, or gases. The potential accuracy of NI to measure b has not been fully realized in past experiments, due to systematic sources of error. We have used a method which eliminates two of the main sources of error to measure the scattering length of silicon with a relative standard uncertainty of 0.005%. The resulting value, $b=4.1507(2)$ fm, is in agreement with the current accepted value, but has an uncertainty five times smaller. [S1050-2947(98)04808-2]

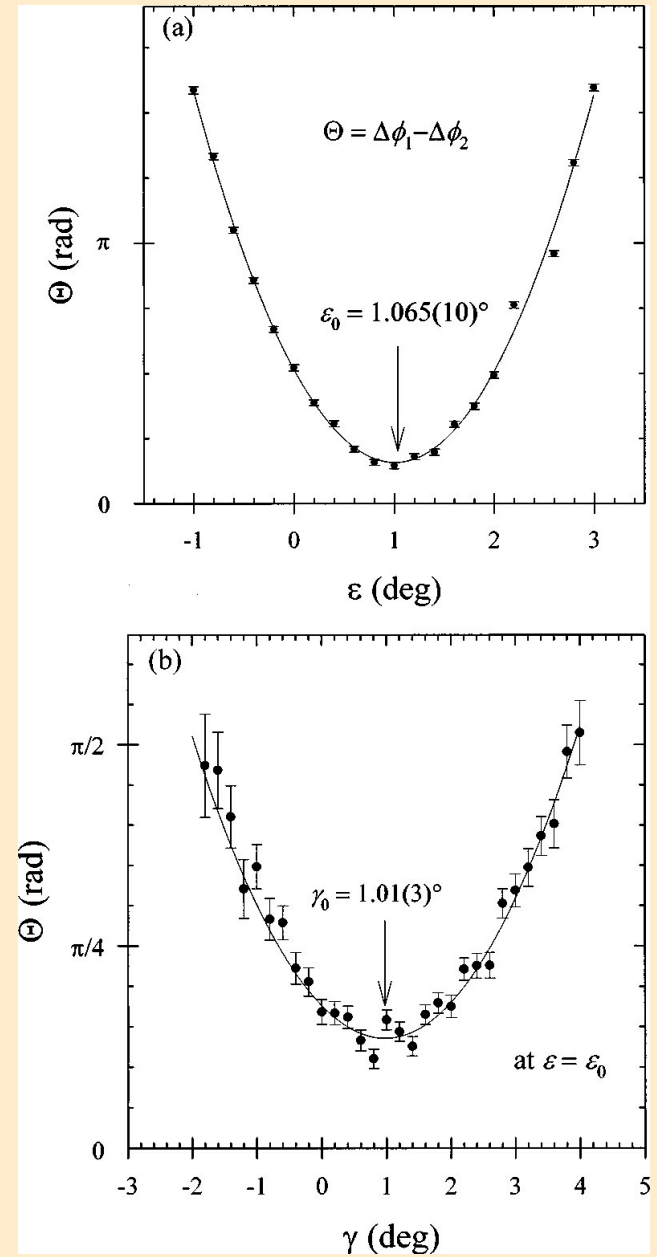
PACS number(s): 03.75.Dg, 07.60.Ly, 61.12.-q



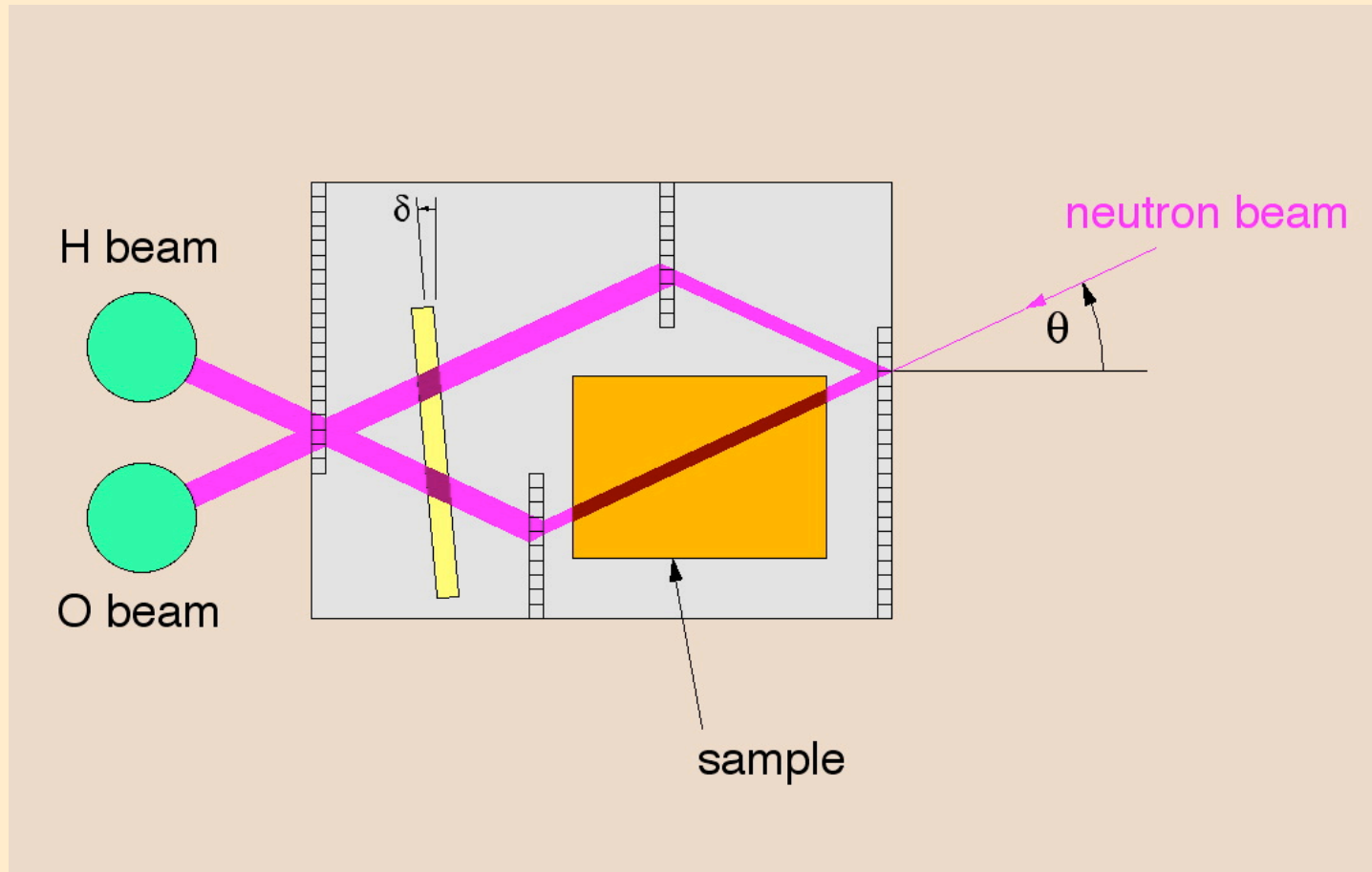


net phase shift: $\Theta(\varepsilon_0, \gamma_0) = 248\pi + 0.455(7)$ radians

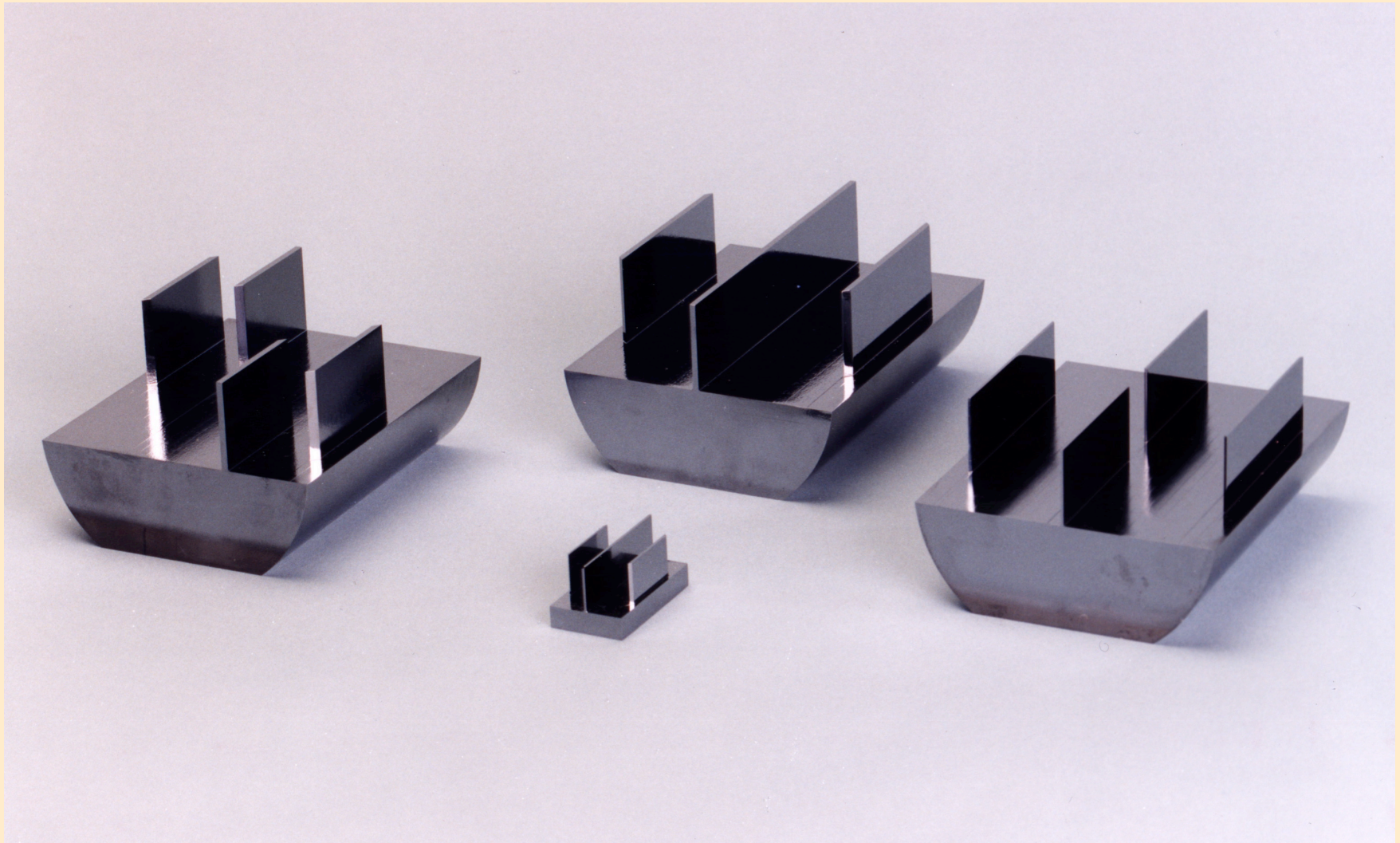
$b_{\text{coh}} = 4.15041(21)$ fm



Skew-Symmetric Neutron Interferometer

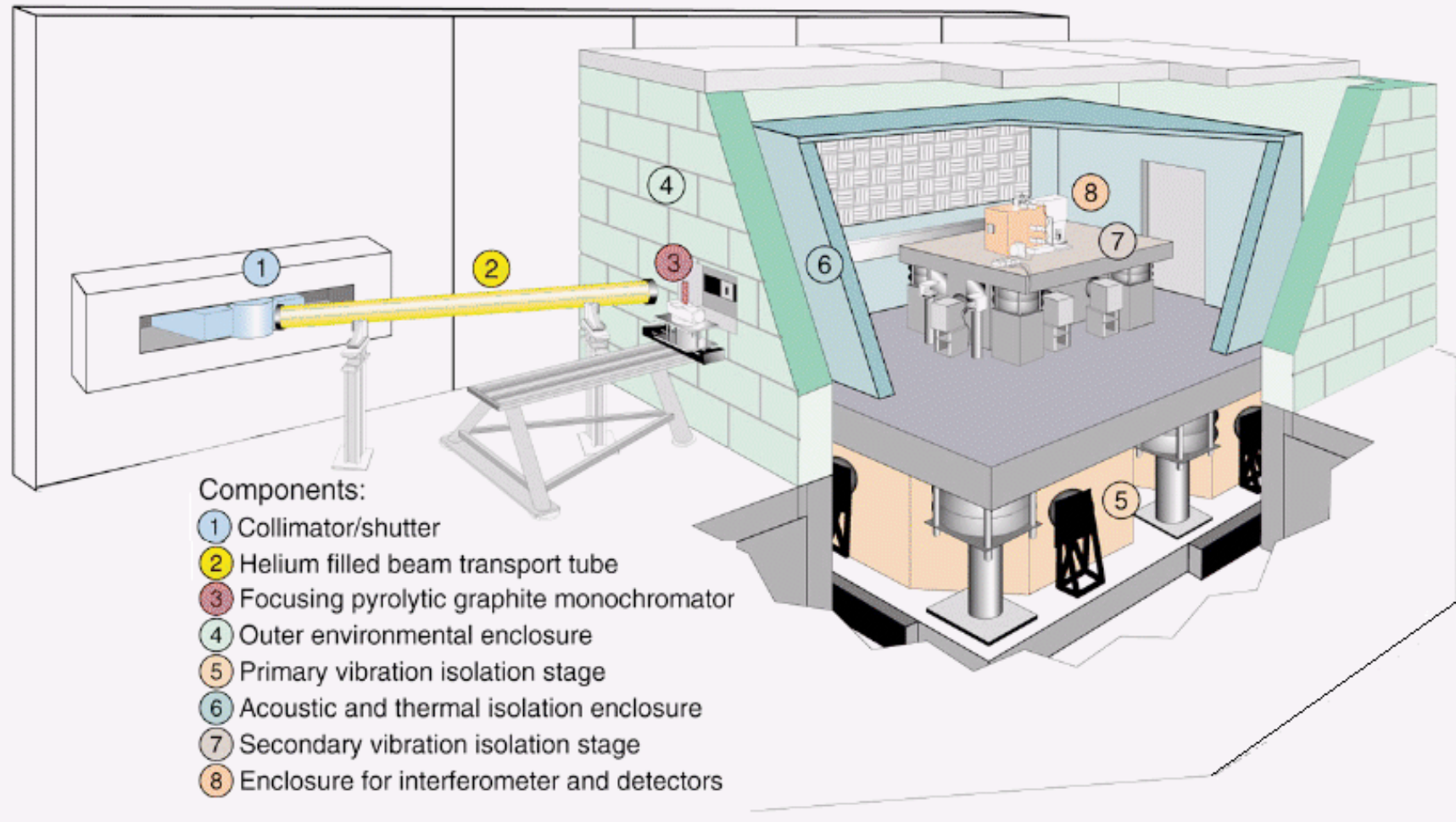


NIST perfect crystal silicon interferometers

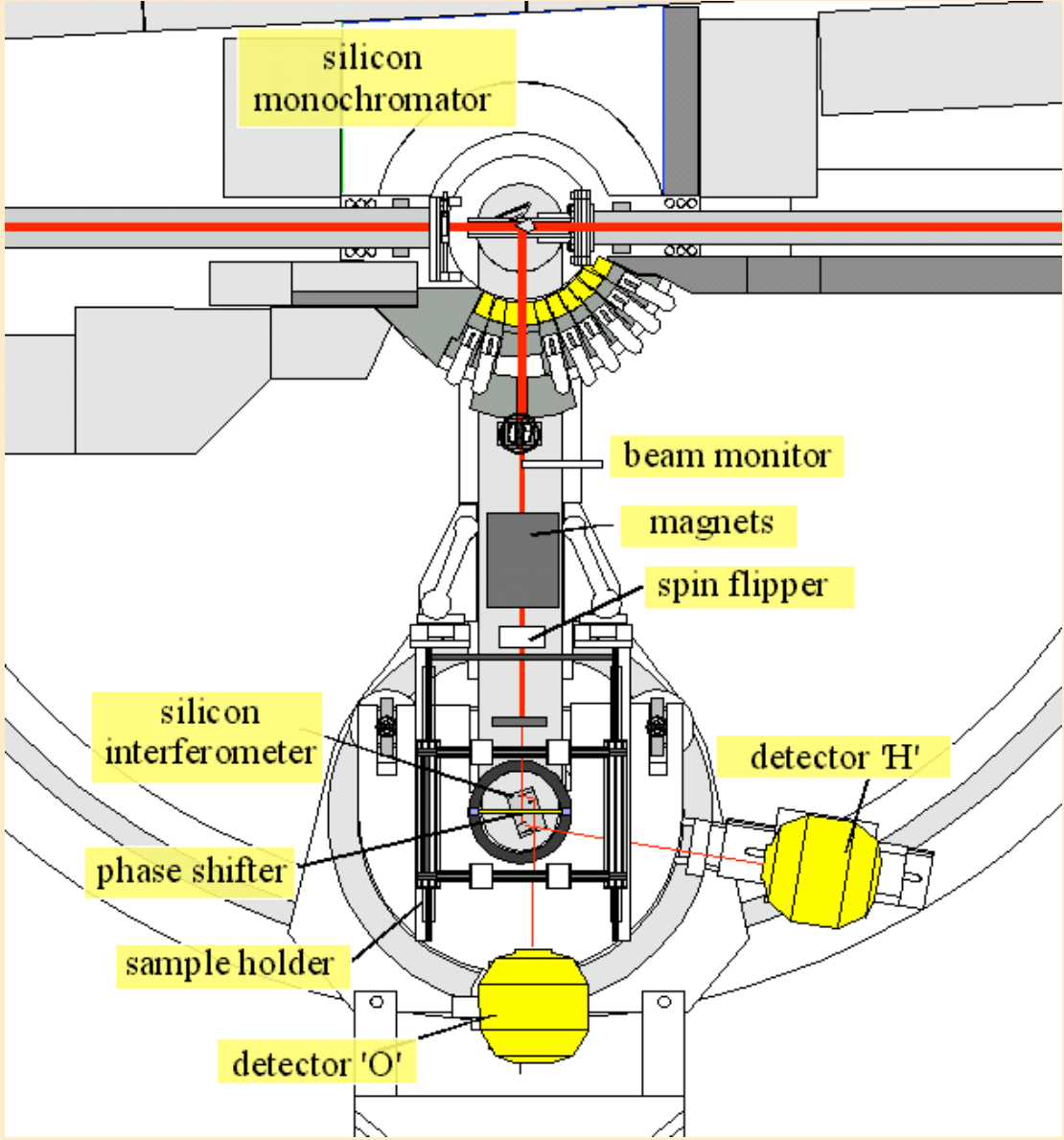


NIST

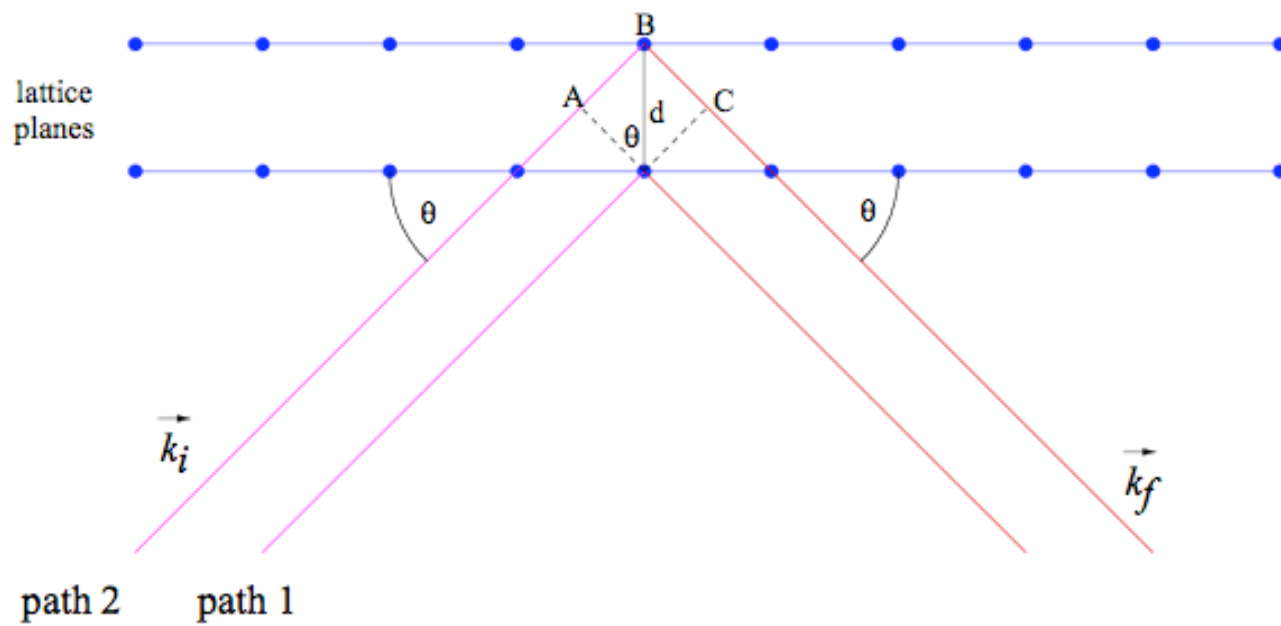
Neutron Interferometer and Optics Facility



S18 Neutron Interferometer at the Institut Laue-Langevin



"Geometric" Bragg Reflection

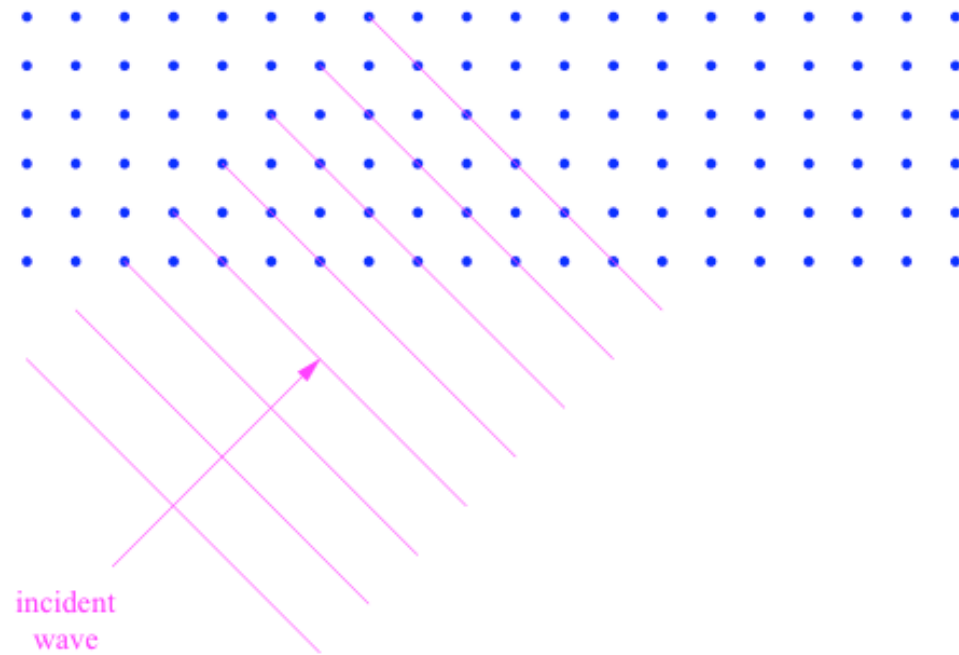


path 2 travels additional distance $\overline{AB} + \overline{BC} = 2d \sin\theta$

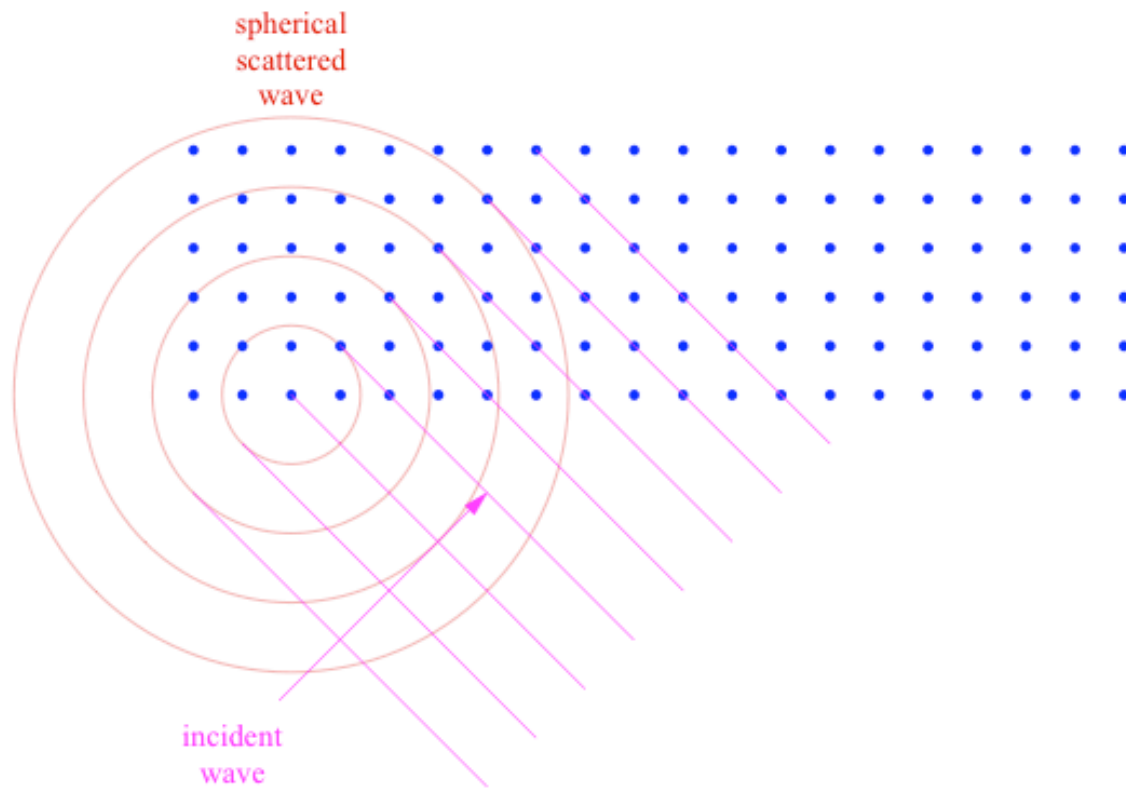
relative phase shift: $\Delta\phi = \frac{2d \sin\theta}{\lambda} \times 2\pi$

condition for constructive interference: $n\lambda = 2d \sin\theta$ (Bragg's Law)

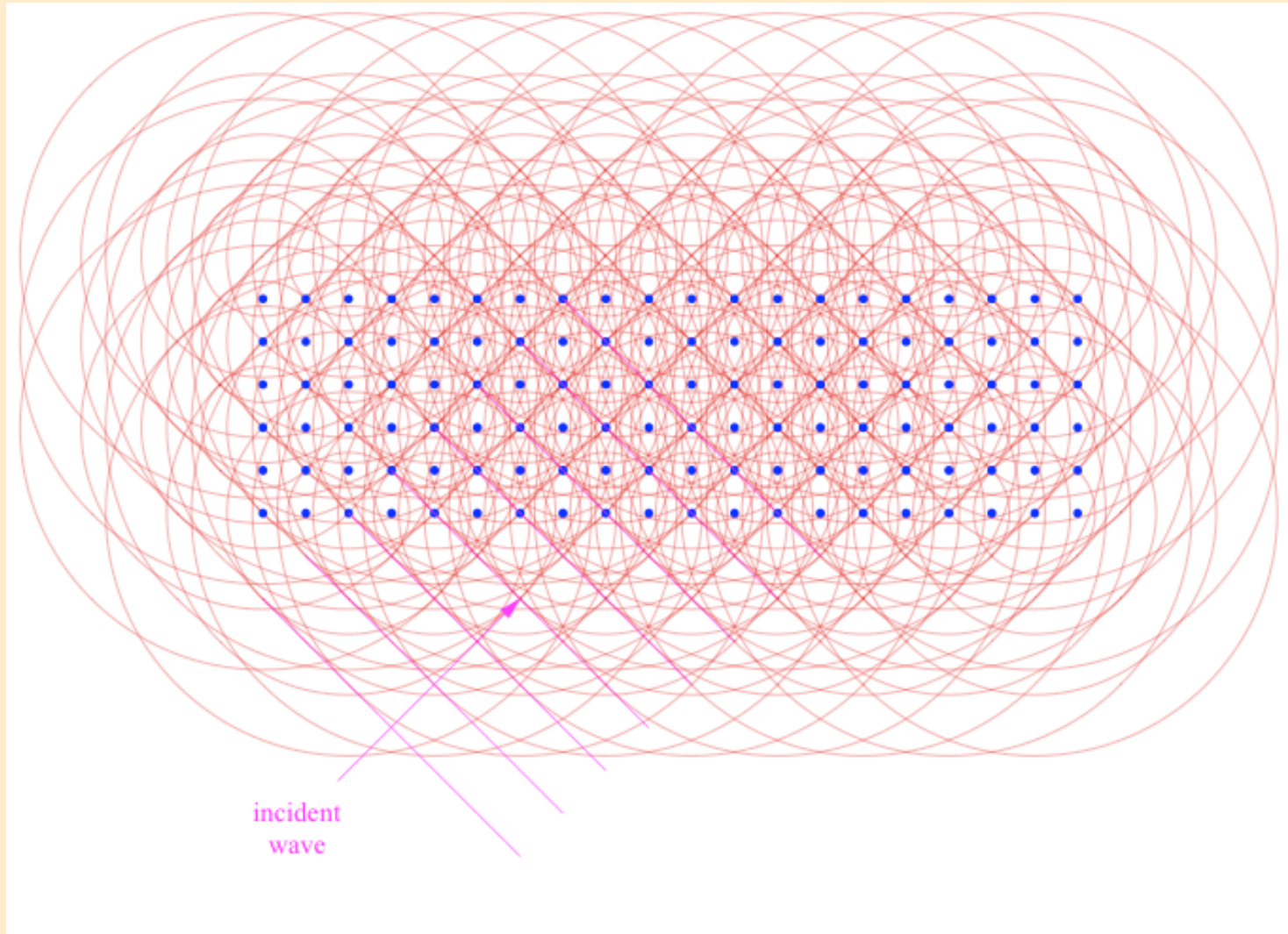
Kinematic Bragg Diffraction



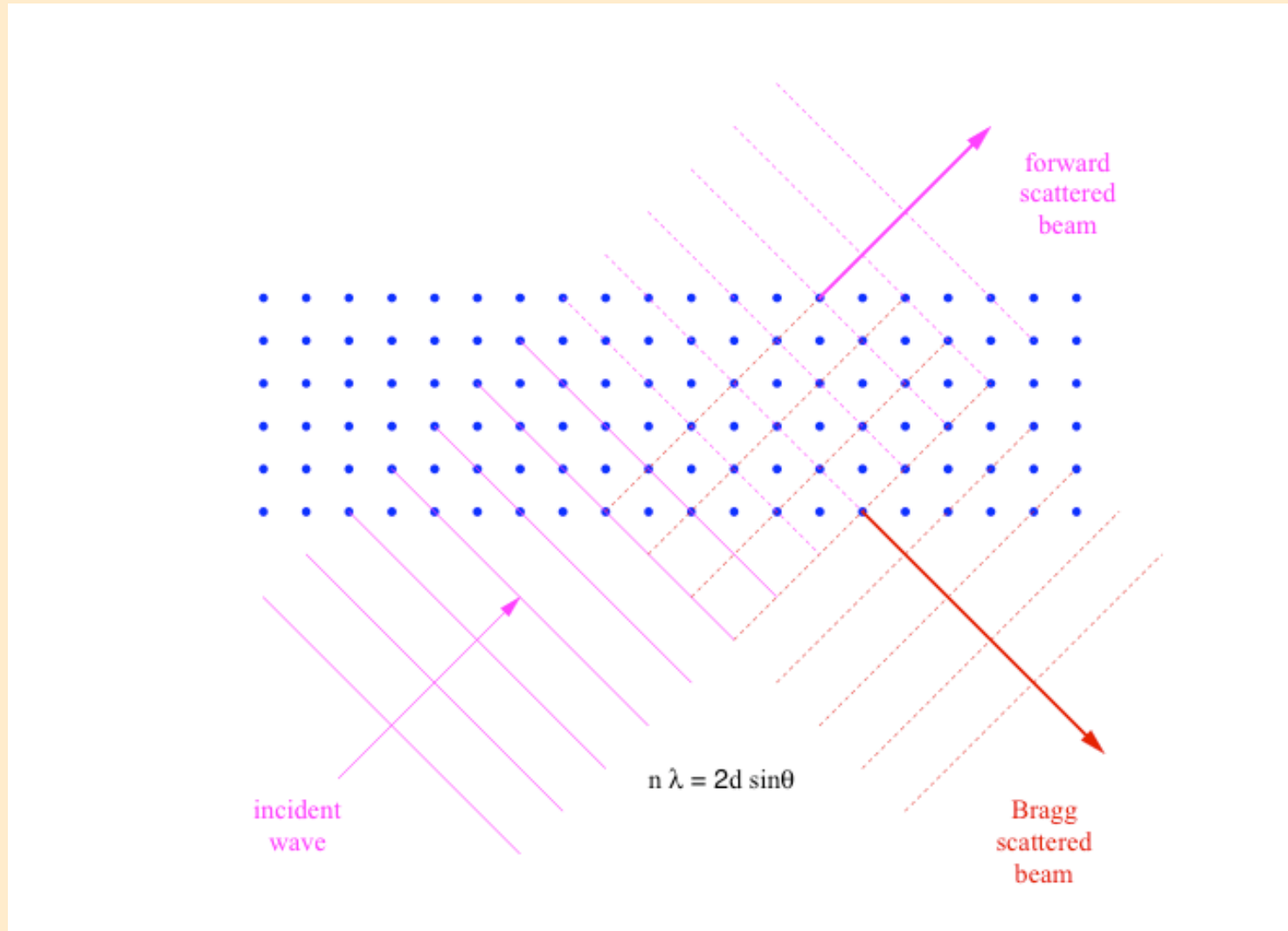
Kinematic Bragg Diffraction



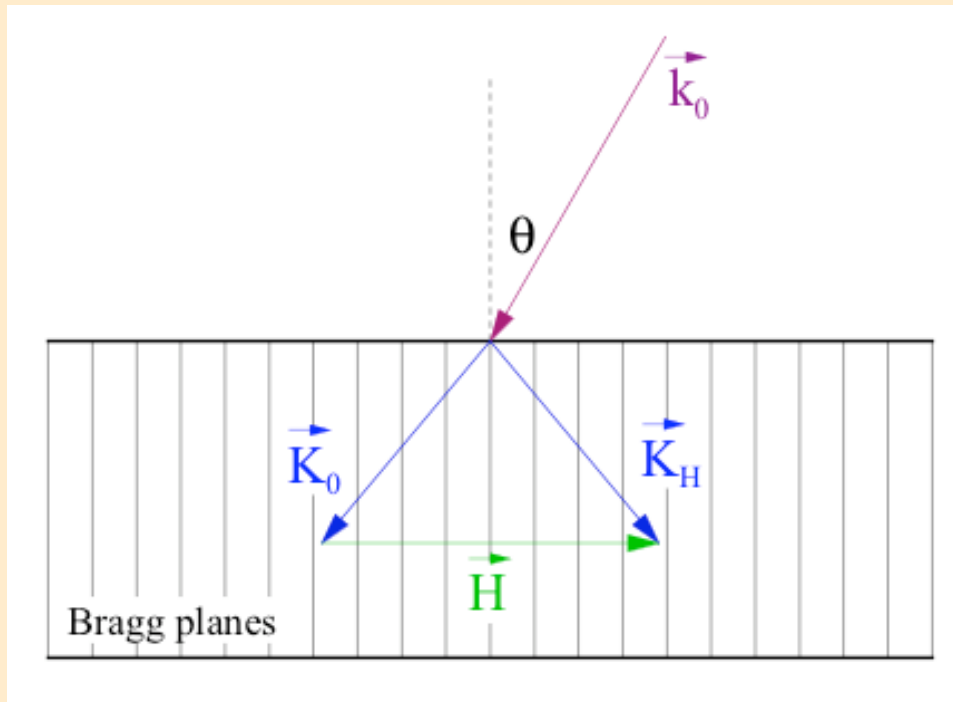
Kinematic Bragg Diffraction



Kinematic Bragg Diffraction



Dynamical Diffraction Theory



\vec{H} = Bragg vector

$$|\vec{H}| = \frac{2\pi n}{d}$$

\vec{K}_0 = internal forward scattered wave

\vec{K}_H = external forward scattered wave

Bragg condition:

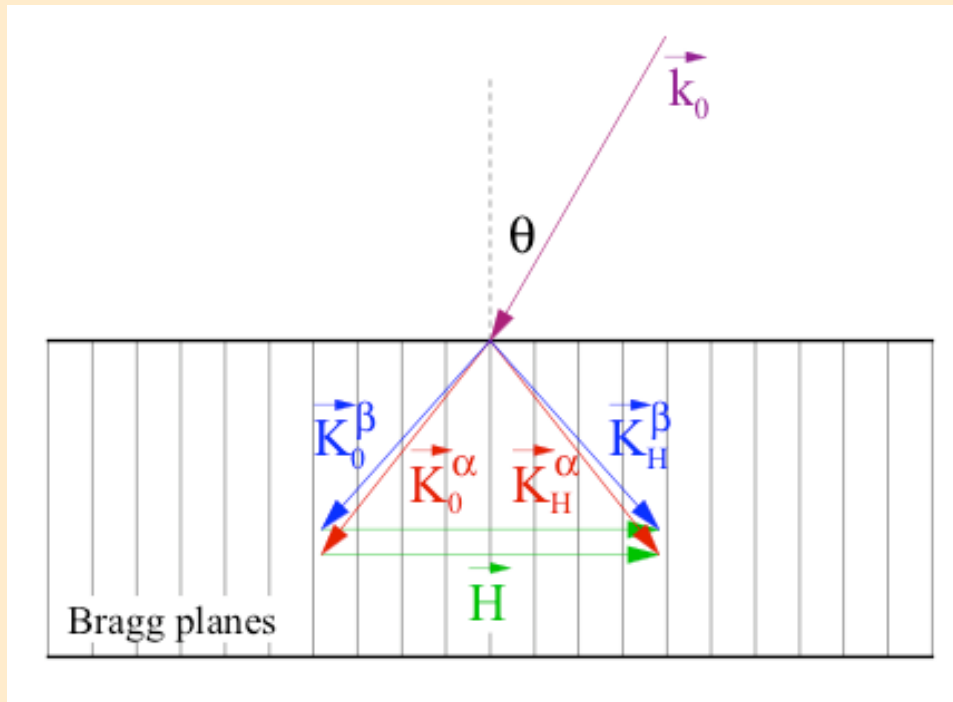
$$\vec{K}_H - \vec{K}_0 = \vec{H}$$

Solve Schrödinger Eqn. inside crystal:

$$\left(\nabla^2 + k_0^2\right)\Psi(\vec{r}) = v(\vec{r})\Psi(\vec{r})$$

$$\text{with } v(\vec{r}) = 4\pi \sum_i b_i \delta(\vec{r} - \vec{r}_i) = \sum_n v_{H_n} e^{i\vec{H}_n \cdot \vec{r}}$$

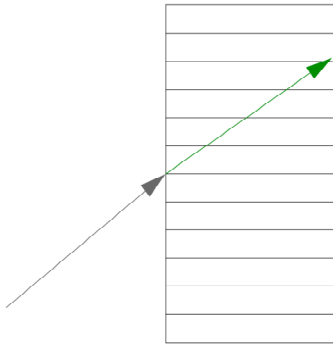
Dynamical Diffraction Theory



Dispersion Equation: $(K^2 - K_0^2)(K^2 - K_H^2) = v_H^2$

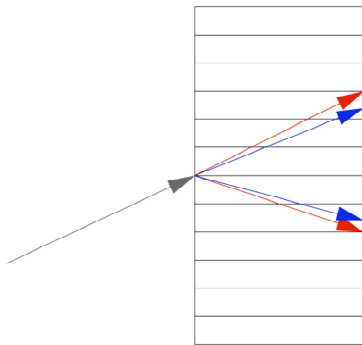
approximate: $(K - K_0)(K - K_H) = \frac{v_H^2}{4k_0^2}$ quadratic equation
2 solutions for K_0

Dynamical Diffraction Theory



off Bragg

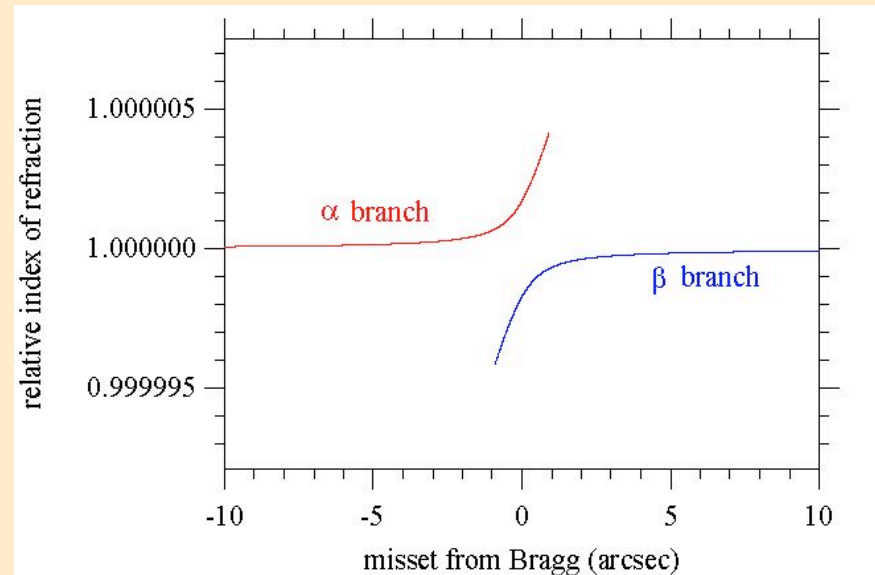
1 solution



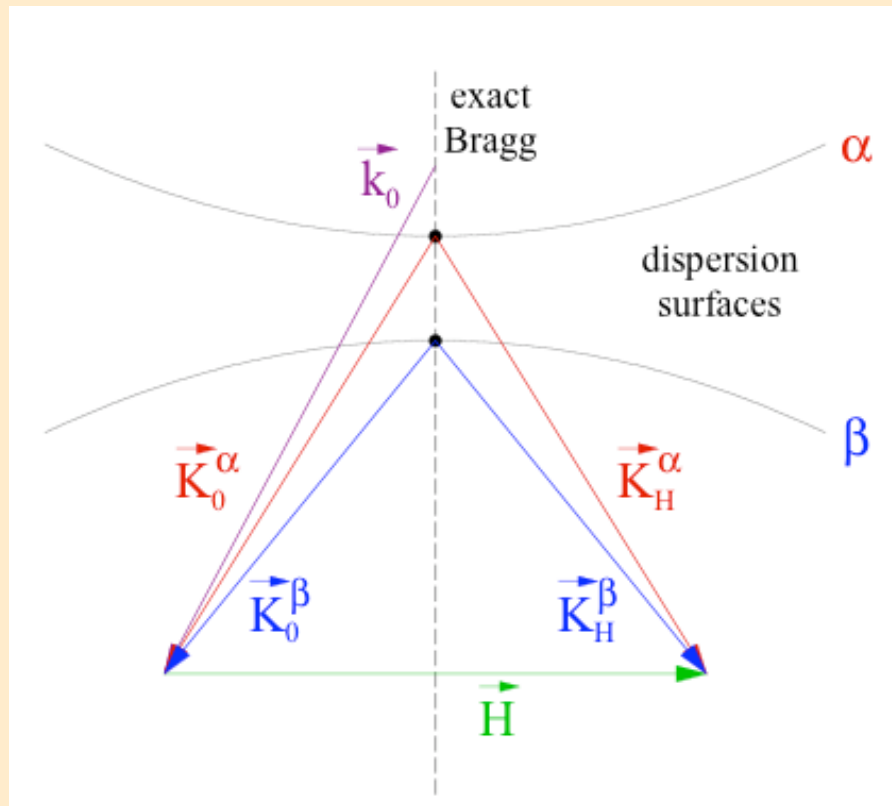
on Bragg

4 solutions

index of refraction
is double-valued

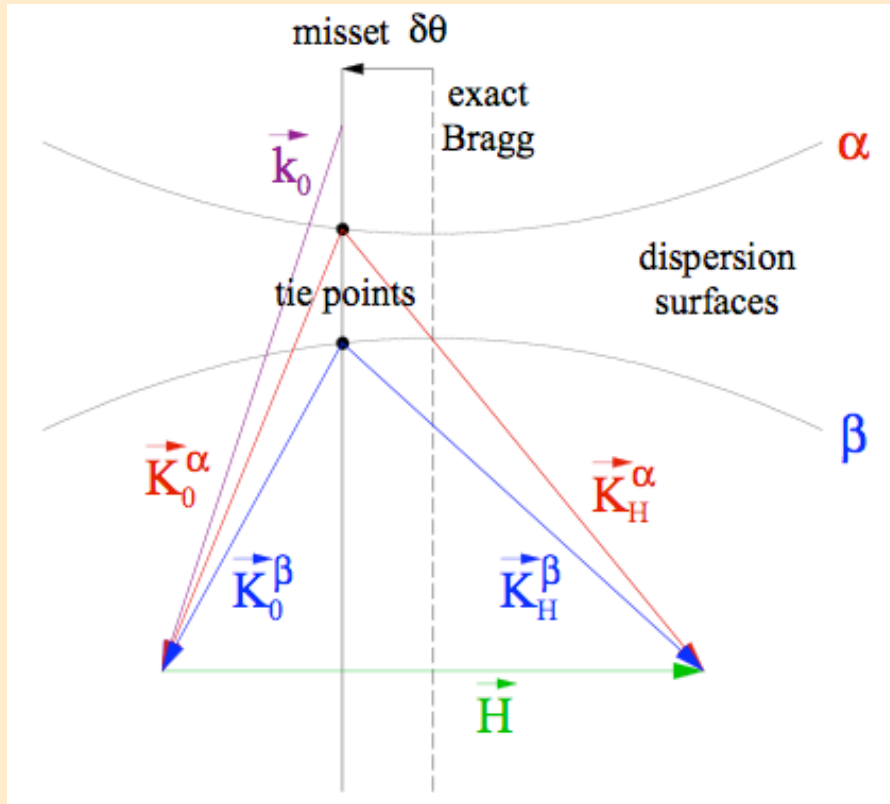


Dynamical Diffraction Theory



internal wave function:
$$\Psi(\vec{r}) = \psi_0^\alpha e^{i\vec{K}_0^\alpha \cdot \vec{r}} + \psi_0^\beta e^{i\vec{K}_0^\beta \cdot \vec{r}} + \psi_H^\alpha e^{i\vec{K}_H^\alpha \cdot \vec{r}} + \psi_H^\beta e^{i\vec{K}_H^\beta \cdot \vec{r}}$$

Dynamical Diffraction Theory



$$\psi_0^\alpha = \frac{1}{2} \left[1 - \frac{y}{\sqrt{1+y^2}} \right] A_0$$

$$\psi_0^\beta = \frac{1}{2} \left[1 + \frac{y}{\sqrt{1+y^2}} \right] A_0$$

$$\psi_H^\alpha = -\frac{1}{2} \left[\frac{1}{\sqrt{1+y^2}} \right] A_0$$

$$\psi_H^\beta = +\frac{1}{2} \left[\frac{1}{\sqrt{1+y^2}} \right] A_0$$

$$y = \frac{k_0 \sin 2\theta_B}{2\nu_H} \delta\theta$$

misset parameter

internal wave function: $\Psi(\vec{r}) = \psi_0^\alpha e^{i\vec{K}_0^\alpha \cdot \vec{r}} + \psi_0^\beta e^{i\vec{K}_0^\beta \cdot \vec{r}} + \psi_H^\alpha e^{i\vec{K}_H^\alpha \cdot \vec{r}} + \psi_H^\beta e^{i\vec{K}_H^\beta \cdot \vec{r}}$

Dynamical Diffraction Theory

Transmitted wave: $\Psi_{\text{trans}}(\vec{r}) = \psi_{\text{tr}0} e^{i\vec{k}_0 \cdot \vec{r}} + \psi_{\text{tr}H} e^{i\vec{k}_H \cdot \vec{r}}$

$$\psi_{\text{tr}0} = \left[\cos \Phi - \frac{iy}{\sqrt{1+y^2}} \sin \Phi \right] e^{i(\phi_1 - \phi_0)} A_0$$

$$\psi_{\text{tr}H} = \left[\frac{-iy}{\sqrt{1+y^2}} \sin \Phi \right] e^{-i(\phi_1 + \phi_0)} A_0$$

with

$$\phi_0 = \frac{v_0 D}{\cos \theta_B}, \quad \phi_1 = \frac{v_H D}{\cos \theta_B}$$

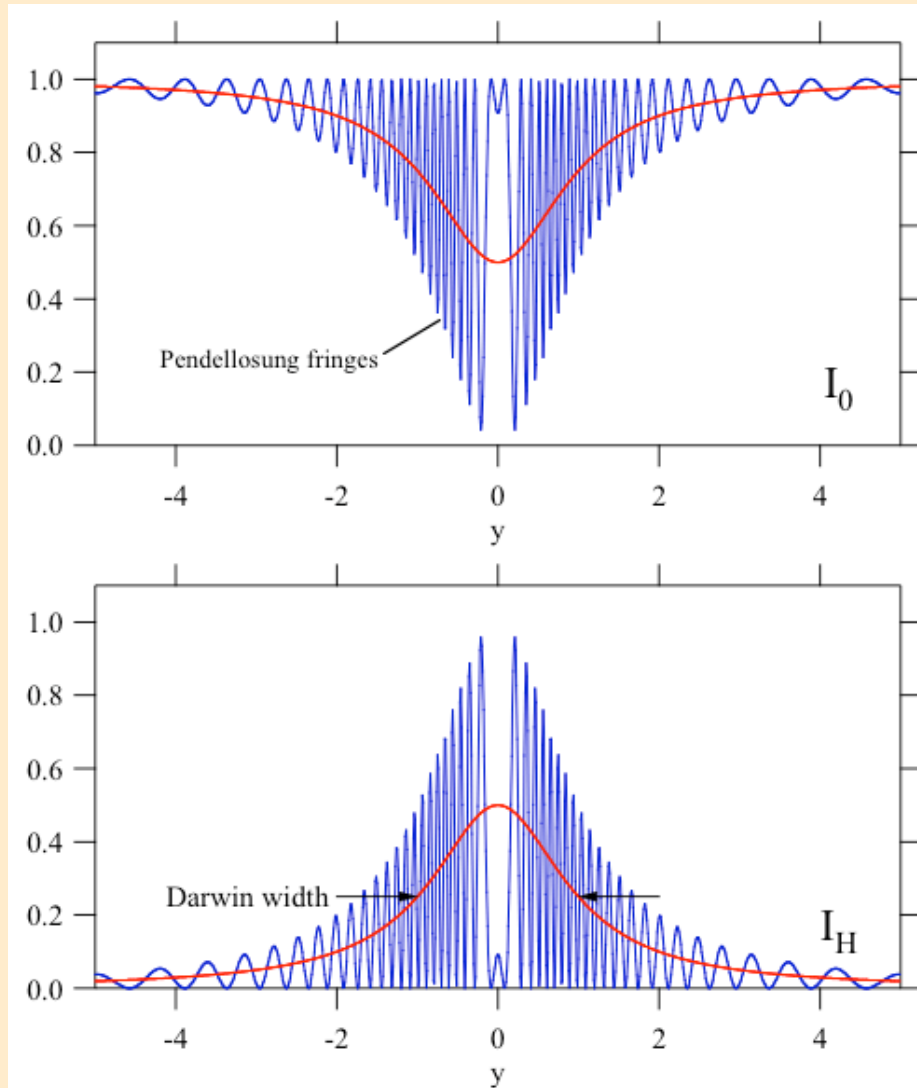
$$\Phi = \left(v_H \frac{1}{\sqrt{1+y^2}} \right) \frac{D}{\cos \theta_B}$$

Transmitted intensities:

$$I_0 = |\psi_{\text{tr}0}|^2 = A_0^2 \left[\cos^2 \Phi + \frac{y^2}{1+y^2} \sin^2 \Phi \right]$$

$$I_H = |\psi_{\text{tr}H}|^2 = A_0^2 \left[\frac{1}{1+y^2} \sin^2 \Phi \right]$$

Transmitted Intensities



For the $\langle 111 \rangle$ reflection in Si
at $\lambda = 2.70 \text{ \AA}$:

$$y = 1 \Rightarrow 0.9 \text{ arcsec}$$

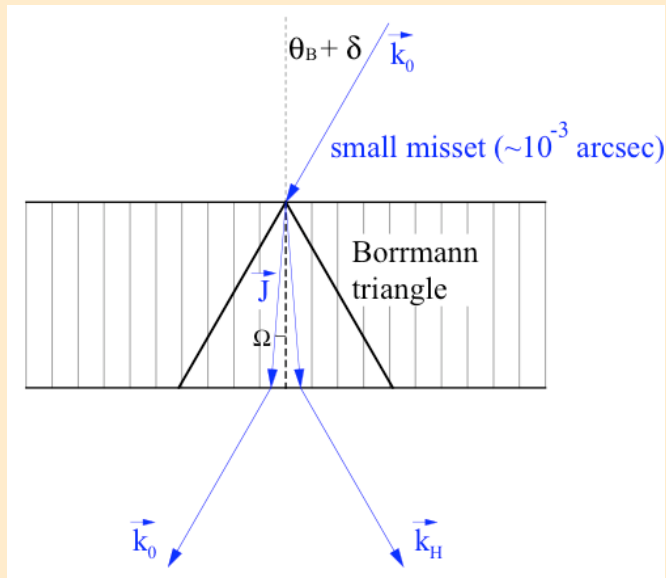
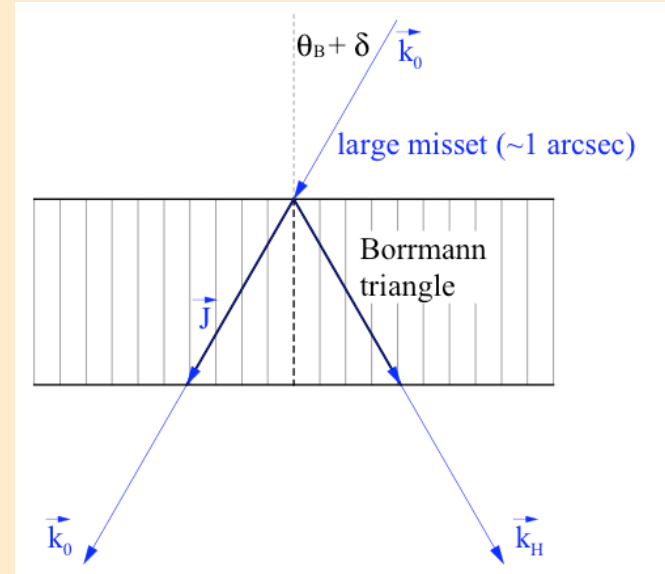
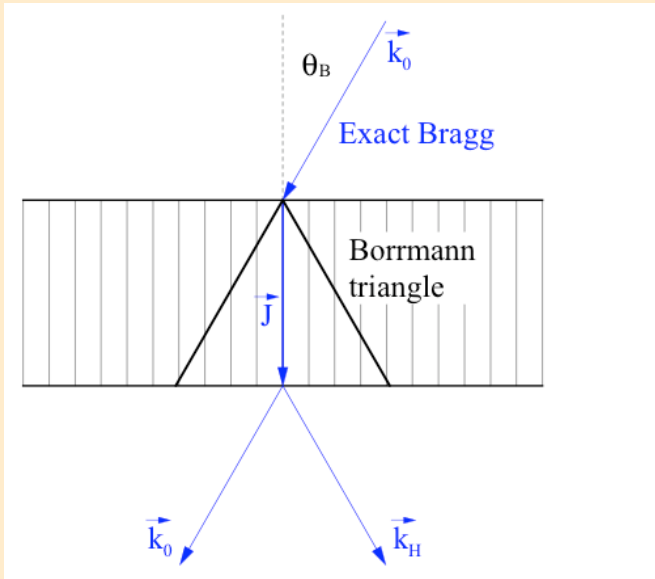
Some Consequences of Dynamical Diffraction

- Pendellösung interference

$$\Phi = \left(v_H \frac{1}{\sqrt{1+y^2}} \right) \frac{D}{\cos \theta_B}$$

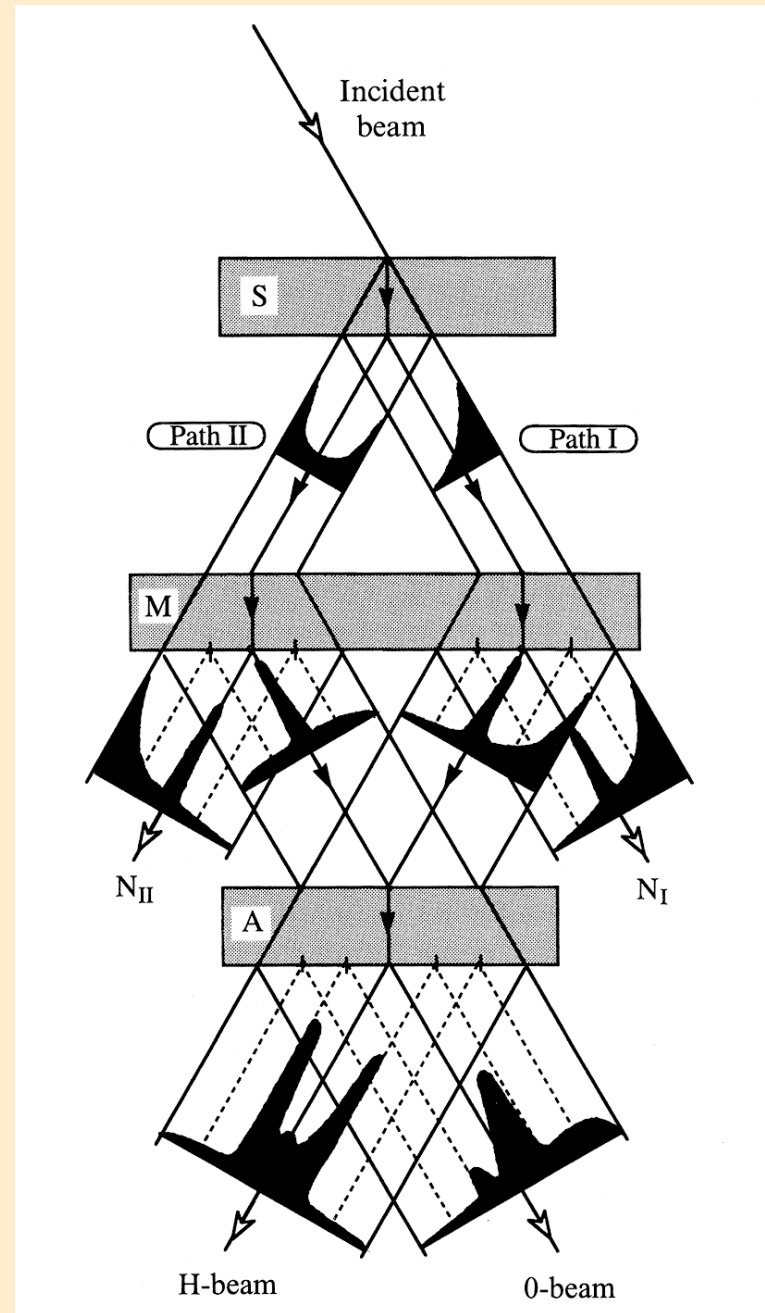
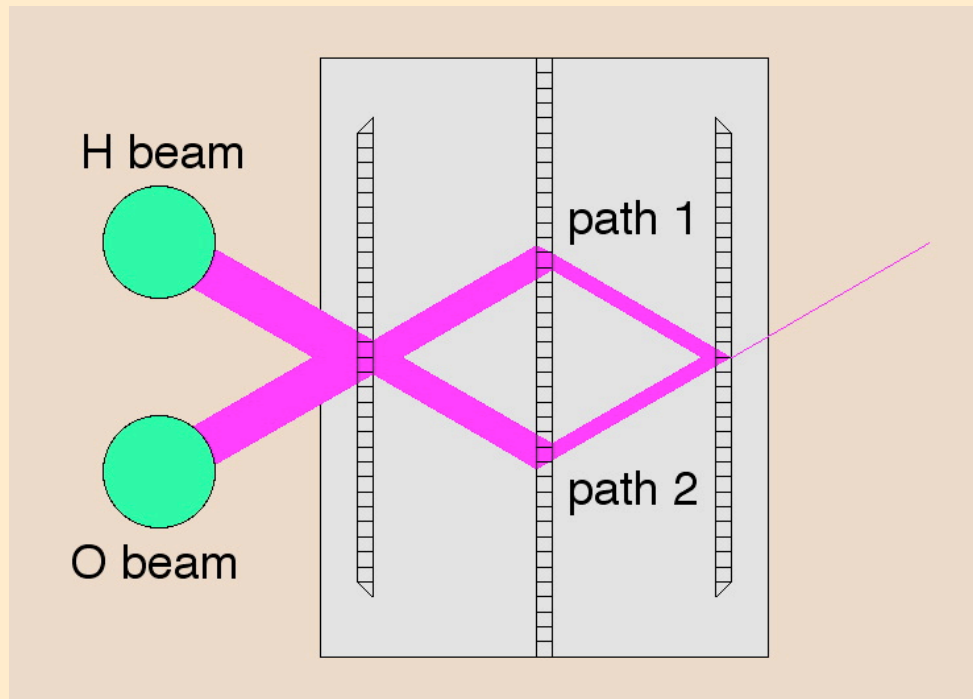
- Anomalous transmission
- Angle amplification

Angle Amplification



For small δ ($\sim 10^{-3}$ arcsec): $\frac{\Omega}{\delta} \approx 10^6$

Practical Neutron Interferometer



4π Rotational Symmetry of Spinors

Rotation operator: $R_{\hat{n}}(\alpha) = e^{-\frac{i}{\hbar}\alpha\hat{n}\cdot\vec{S}}$

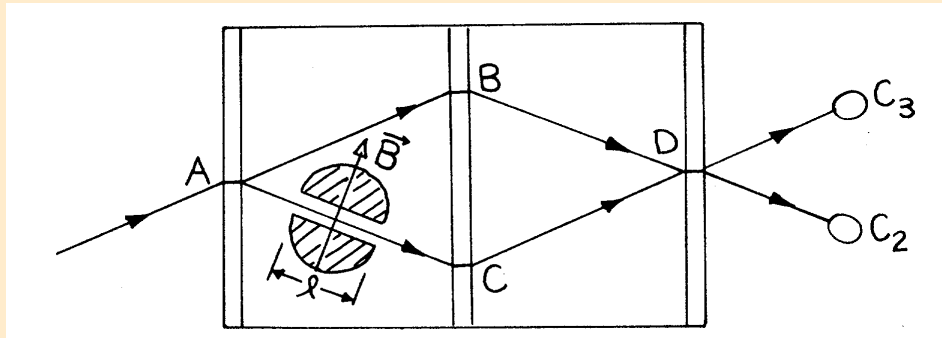
Spin-1/2 particle: $\vec{S} = \frac{1}{2}\hbar\vec{\sigma}$ so $R_{\hat{n}}(\alpha) = e^{-i\frac{\alpha}{2}\hat{n}\cdot\vec{\sigma}}$

Rotations about z-axis: $R_z(\alpha) = \begin{pmatrix} e^{-i\alpha/2} & 0 \\ 0 & e^{i\alpha/2} \end{pmatrix}$

Symmetry:

$$R_z(2\pi)\chi = -\chi$$

$$R_z(4\pi)\chi = \chi$$



Larmor precession phase:

$$\Delta\phi = \pm 2\pi\mu_n m_n \lambda B \ell / \hbar^2$$



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**NUCLEAR
INSTRUMENTS
& METHODS
IN PHYSICS
RESEARCH**
Section A

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4π -Periodicity of the spinor wave function under space rotation

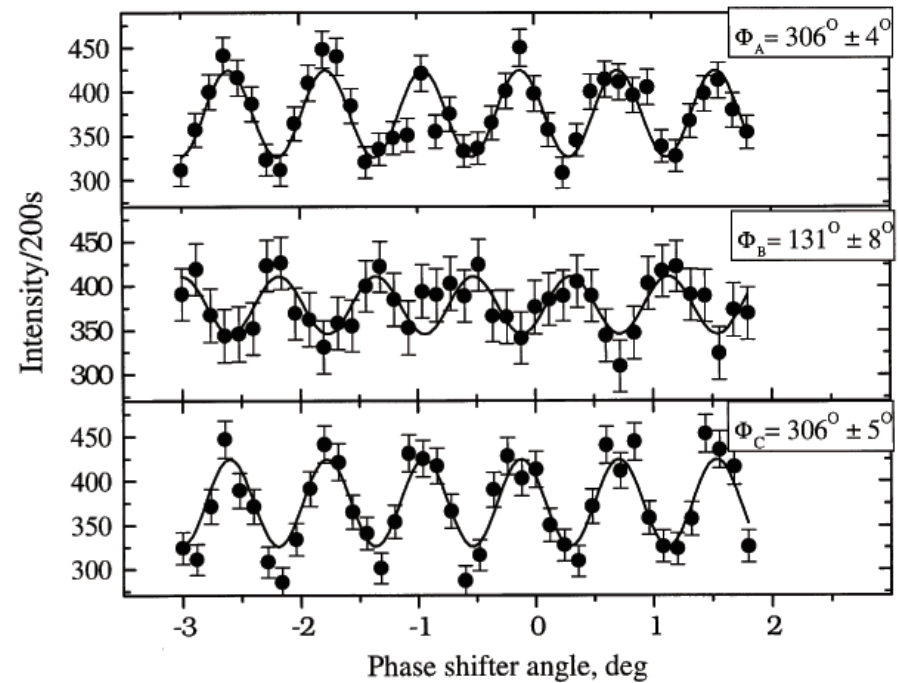
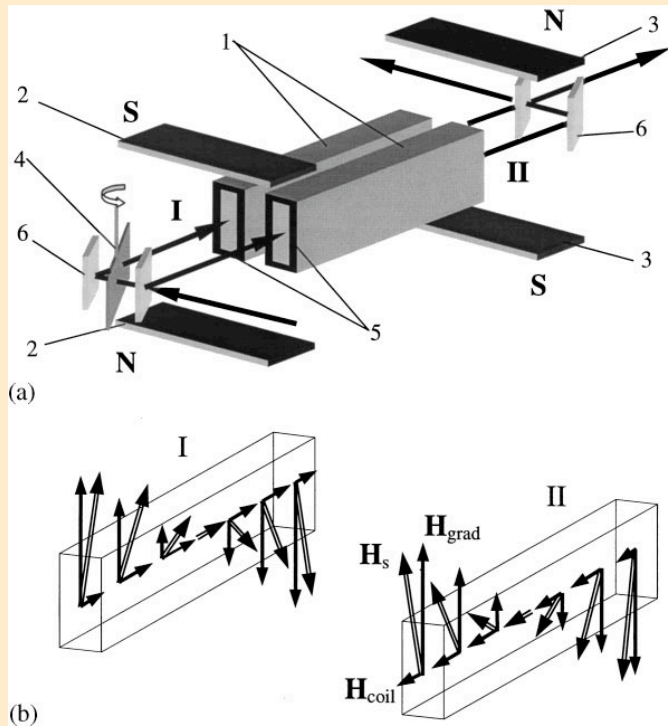
P. Fischer^a, A. Ioffe^{b,c,*}, D.L. Jacobson^c, M. Arif^c, F. Mezei^{a,d}

^aBerlin Neutron Scattering Center, Hahn-Meitner-Institut, Glienicke Str. 100, 14109 Berlin, Germany

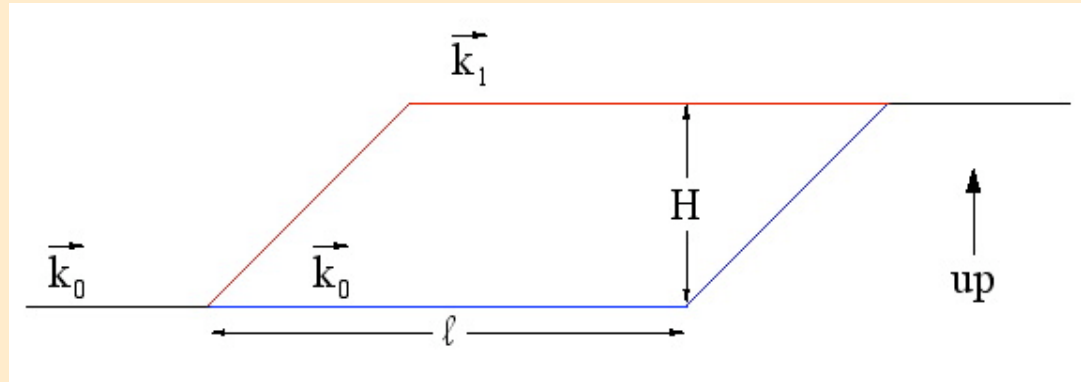
^bDepartment of Physics and Astronomy, University of Missouri-Columbia, Columbia, MO 65211, USA

^cNational Institute of Standards and Technology, Gaithersburg, MD 20899, USA

^dLos Alamos National Laboratory, Los Alamos, NM 87545, USA



Quantum Phase Shift Due To Gravity (COW Experiments)



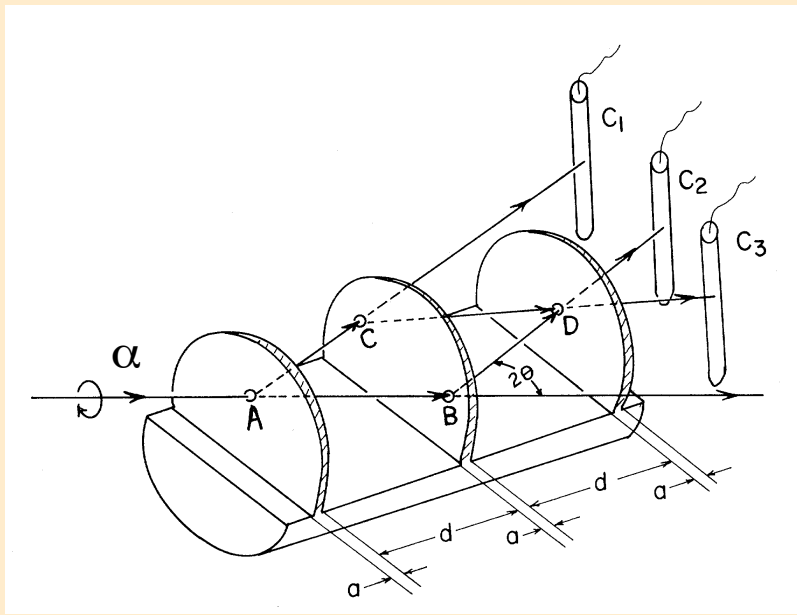
$$\Delta\phi = \frac{2\pi\lambda g A}{h^2} m_{\text{in}} m_{\text{grav}}$$

m_{in} = neutron inertial mass

$A = H\ell$ = area of parallelogram

m_{grav} = neutron gravitational mass

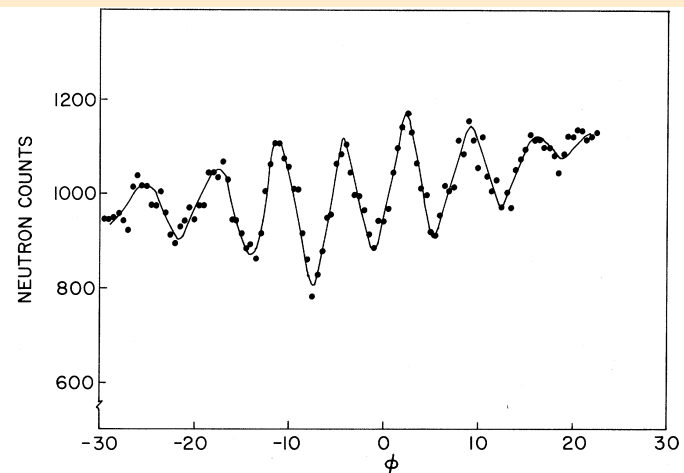
test of weak equivalence principle at the quantum limit



measured: $q = 54.3$
 theory: $q = 59.6$

$$\Delta\phi_{\text{grav}} = \frac{2\pi\lambda g A_0}{h^2} m_{\text{in}} m_{\text{grav}} \sin\alpha = q \sin\alpha$$

$A_0 =$ area of parallelogram at $\alpha = 0$



Systematic Effects in the COW Experiments

$$q_{\text{COW}} = \left[\left(q_{\text{grav}} (1 + \varepsilon) + q_{\text{bend}} \right)^2 + q_{\text{Sagnac}}^2 \right]^{\frac{1}{2}}$$

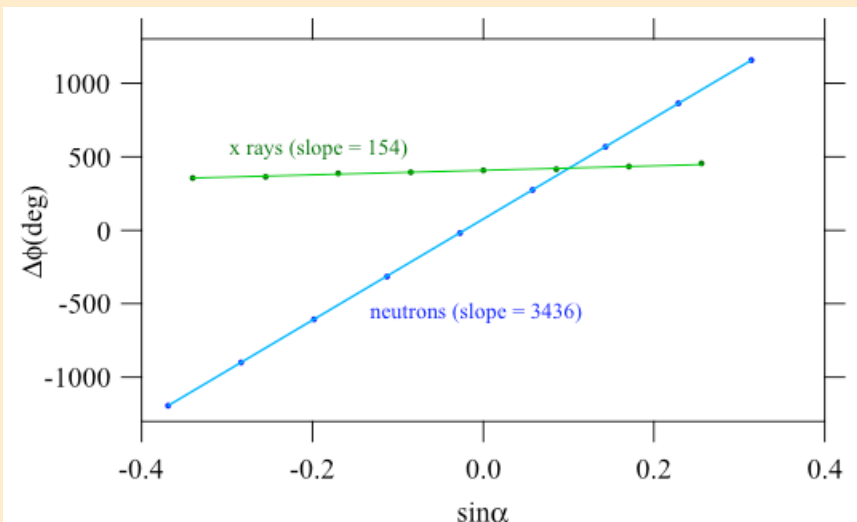
dynamical
diffraction
correction

bending of
interferometer

Earth's rotation

Sagnac effect: $\Delta\phi_{\text{Sagnac}} = \frac{2m_{\text{in}}}{\hbar} \vec{\Omega} \cdot \vec{A}$ due to Earth's rotating frame

bending effect: repeat experiment with x rays, different wavelengths



data from Werner, *et al.* (1988)

Layer and Greene (1991): x rays do not fill the Borrmann fan as completely as neutrons

Littrell, *et al.* (1997) results:

experiment	q_{COW} theory [rad]	q_{COW} meas. [rad]	discrepancy (%)
SS, 440	50.97(5)	50.18(5)	-1.6
SS, 220	100.57(10)	99.02(10)	-1.5
LLL, 440	113.60(10)	112.62(15)	-0.9
LLL, 220	223.80(10)	221.85(30)	-0.9

Upcoming new effort (H. Kaiser, S. Werner, FEW, et al.):
Suspend interferometer inside chamber filled with $\text{ZnBr}_2 + \text{D}_2\text{O}$ (floating COW)

Measuring the Neutron's Mean Square Charge Radius Using Neutron Interferometry

F. E. Wietfeldt, M. Huber

Tulane University, New Orleans, USA

M. Arif, D. L. Jacobson, S. A. Werner

National Institute of Standards and Technology, Gaithersburg, USA

T. C. Black

University of North Carolina, Wilmington, USA

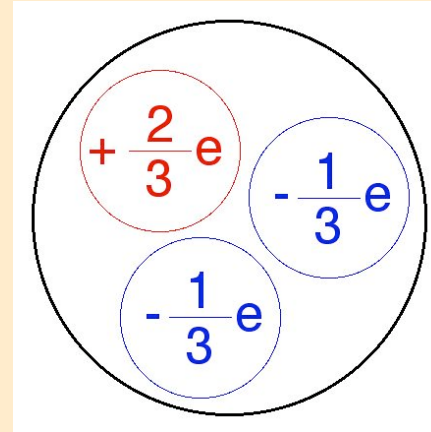
H. Kaiser

Indiana University, Bloomington, USA

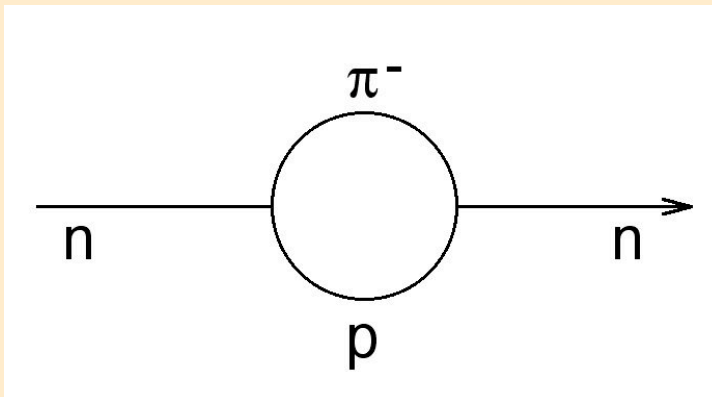
neutron: neutral but consists of charged quarks

neutron mean square charge radius:

$$\langle r_n^2 \rangle = \int \rho(r) r^2 d^3r$$



expected to be negative (positive core, negative skin):



Fermi and Marshall, 1947

Neutron Electric Scattering Form Factor

$G_E^n(Q^2)$ = Fourier transform of neutron charge density (Breit frame)

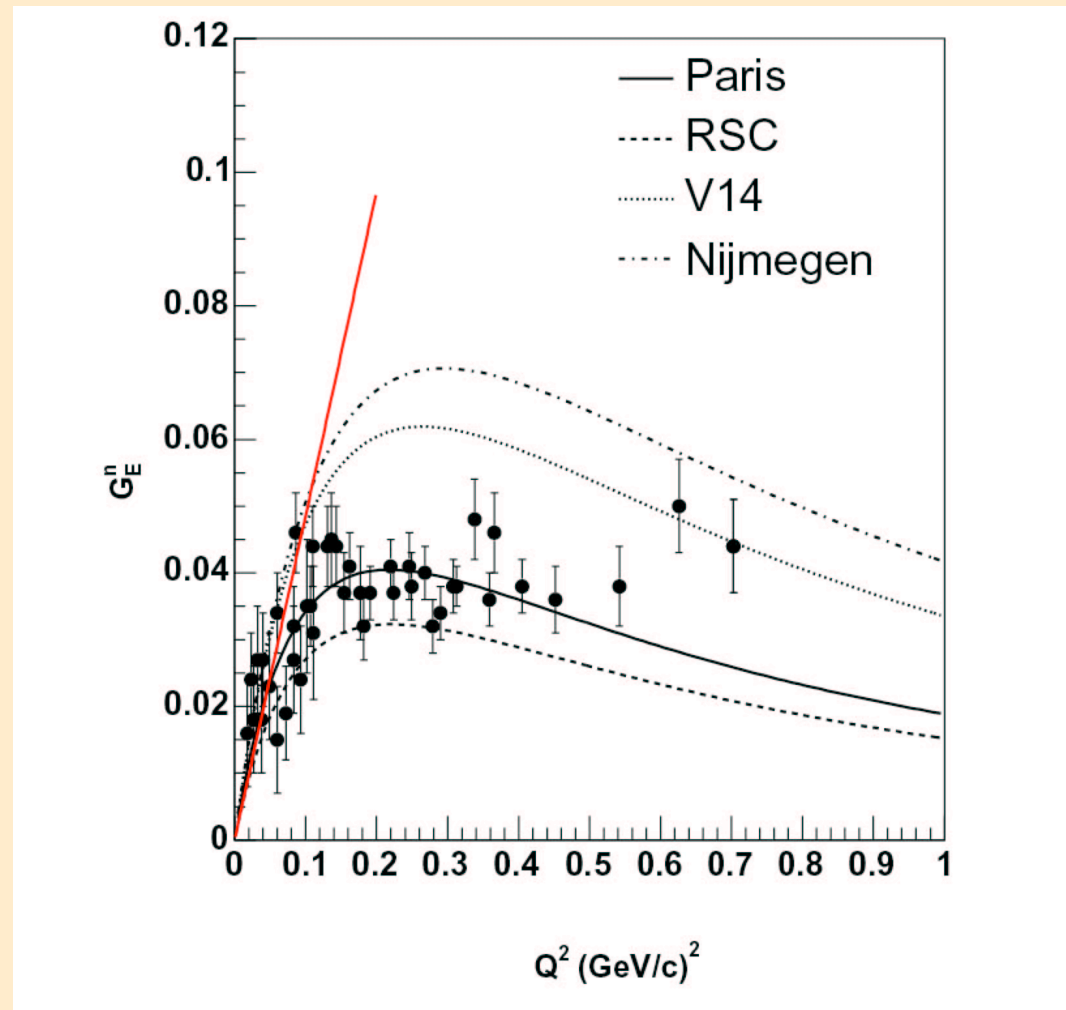
Expanding in momentum transfer Q^2 :

$$G_E^n(Q^2) = q_n - \frac{1}{6} \langle r_n^2 \rangle Q^2 + \dots$$

In the low Q^2 limit:

$$\langle r_n^2 \rangle = -6 \left. \frac{dG_E(Q^2)}{dQ^2} \right|_{Q^2=0}$$

$\langle r_n^2 \rangle$ constrains the slope of $G_E(Q^2)$ in electron scattering experiments and theory (e.g. Bates, Jefferson Lab)



Neutron-Atom Coherent Scattering Length

$$b_{\text{coh}} = b_N + Z[1 - f(q)]b_{ne}$$

Fourier transform of
charge density

$$f(q) = \frac{1}{\sqrt{2\pi}} \int e^{iq \cdot r} \rho_{\text{atom}}(r) d^3r$$

b_{ne} = neutron-electron scattering length

In 1st Born approximation: $\langle r_n^2 \rangle = 3a_0 \left(\frac{m_e}{m_n} \right) b_{ne} = (86.34 \text{ fm}) b_{ne}$

Foldy Scattering Length

$$b_F = -\frac{\gamma e^2}{2m_e c^2} = -1.468 \times 10^{-3} \text{ fm} \quad \text{from neutron's magnetic moment}$$

Incorrect interpretation: $b_{ne} (\text{meas.}) = b_{\text{intrinsic}} + b_F$

Correct interpretation: The experimentally measured value of b_{ne} is *entirely* due to the static charge distribution in the neutron.

[N. Isgur, Phys. Rev. Lett. **83**, 272 (1999)]

Previous Experiments

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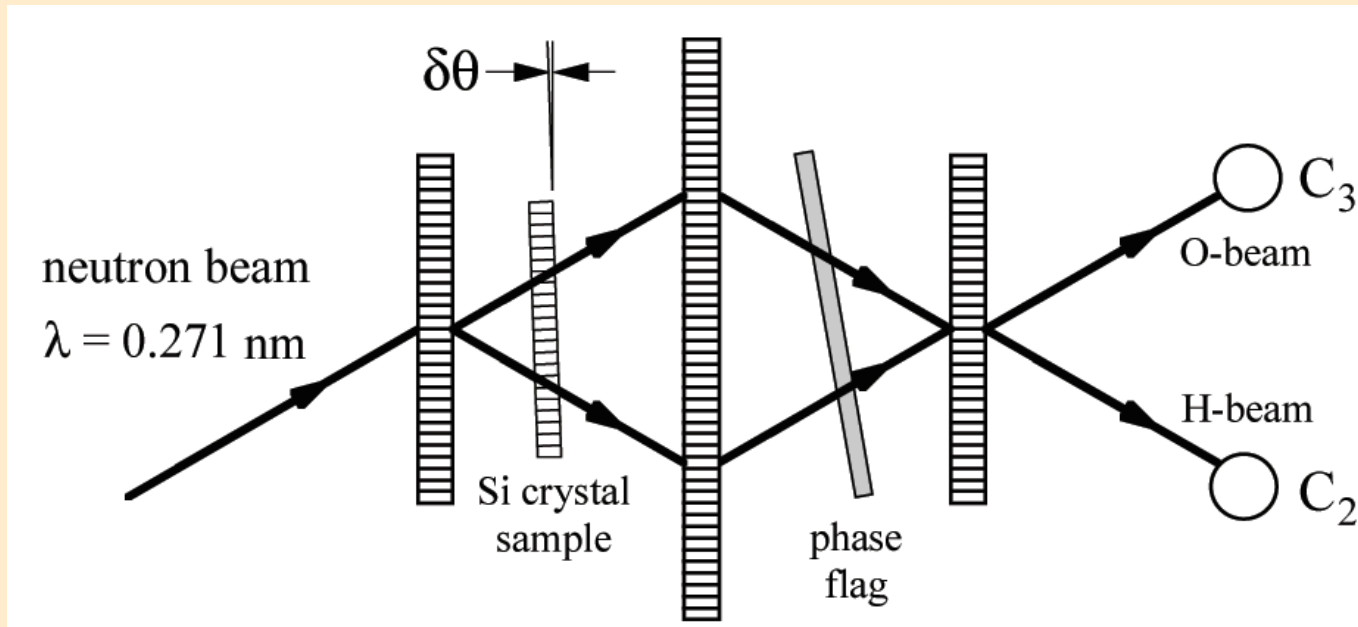
S. KOPECKY *et al.*

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TABLE I. Experimental results of b_{ne} in units of 10^{-3} fm.

Experiment	Target	Result	Reference
Angular scattering	Ar	-0.1 ± 1.8	1947 [7] Fermi
Transmission	Bi	-1.9 ± 0.4	1951 [8] Havens
Angular scattering	Kr, Xe	-1.5 ± 0.4	1952 [9] Hamermesh
Mirror reflection	Bi/O	-1.39 ± 0.13	1953 [10] Hughes
Angular scattering	Kr, Xe	-1.4 ± 0.3	1956 [11] Crouch
Crystal spectrometer transmission	Bi	-1.56 ± 0.05	1959 [2] Melkonian
		-1.49 ± 0.05	1976 in Ref. [15]
		$-1.44 \pm 0.033 \pm 0.06$	1997 this work
Angular scattering	Ne, Ar, Kr, Xe	-1.34 ± 0.03	1966 [12] Krohn
Angular scattering	Ne, Ar, Kr, Xe	-1.30 ± 0.03	1973 [13] Krohn
Single crystal scattering	^{186}W	-1.60 ± 0.05	1975 [14] Alexandrov
Filter-transmission, mirror reflection	Pb	-1.364 ± 0.025	1976 [15] Koester
Filter-transmission, mirror reflection	Bi	-1.393 ± 0.025	1976 [15] Koester
n -TOF transmission, mirror reflection Ref. [17]	Bi	-1.55 ± 0.11	1986 [16] Alexandrov
Filter-transmission, mirror reflection	Pb, Bi	-1.32 ± 0.04	1986 [17] Koester
n -TOF transmission	thorogenic ^{208}Pb	$-1.31 \pm 0.03 \pm 0.04$	1995 [1] Kopecky
		$-1.33 \pm 0.027 \pm 0.03$	1997 this work
Filter-transmission, mirror reflection	Pb-isotopes, Bi	-1.32 ± 0.03	1995 [5] Koester
Garching-Argonne compilation	[12,13,15,17]	-1.31 ± 0.03	1986 [3] Sears
Dubna compilation	[14,16]	-1.59 ± 0.04	1989 [19] Alexandrov
Foldy approximation, b_F		-1.468	1952 [18] Foldy

Neutron Interferometer Experiment



off Bragg: $b_{\text{coh}} = b_N + Z[1 - f(0)]b_{ne} = b_N$

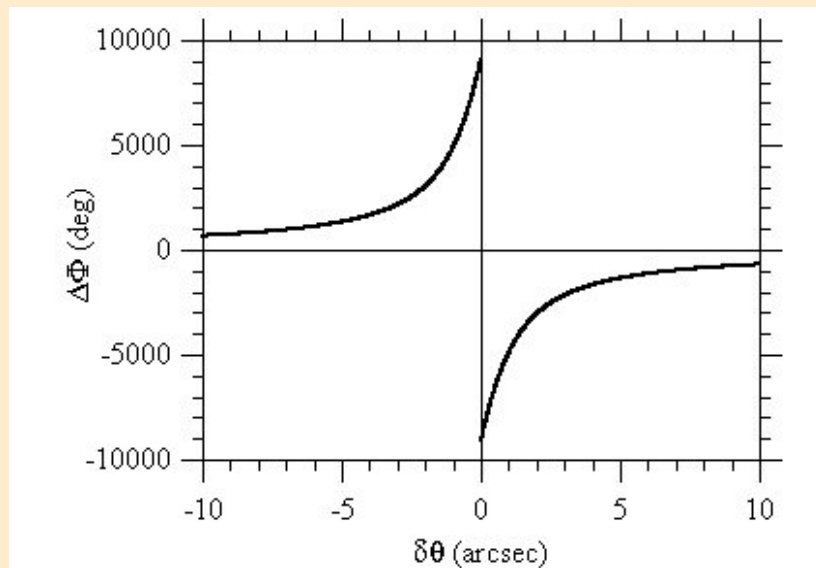
near Bragg: $b_{\text{coh}} = b_N + Z[1 - f(\vec{H}_{111})]b_{ne}$

Dynamical Phase Shift Through Bragg

$$\Delta\Phi_{\text{dyn}} = \frac{v_H}{\cos\theta_B} \left(y \pm \sqrt{1 + y^2} \right) D$$

D = crystal thickness

$$\text{scaled misset angle } y = \frac{k \sin 2\theta_B}{2v_H}$$



$$v_H = \frac{F_{111}\lambda}{V_{\text{cell}}} = \frac{\sqrt{32}\lambda}{V_{\text{cell}}} b_{\text{coh}}$$

$$\text{near Bragg: } b_{\text{coh}} = b_N + Z[1 - f(\vec{H}_{111})]b_{ne}$$

What we must measure:

1. Net dynamical phase shift through Bragg $\rightarrow v_H \rightarrow b_N + Z[1 - f(\vec{H}_{111})]b_{ne}$
to $\sim 10^{-5}$

The maximum slope is $\sim 88\pi/\text{arcsec}$ so we need 0.01 arcsec angular precision to detect every 2π of phase shift

2. Forward phase shift off Bragg $\rightarrow b_N$ to $\sim 10^{-5}$ and subtract
3. Neutron wavelength to $\sim 10^{-3}$
4. Calculate $f(\vec{H}_{111})$ to $\sim 10^{-3}$

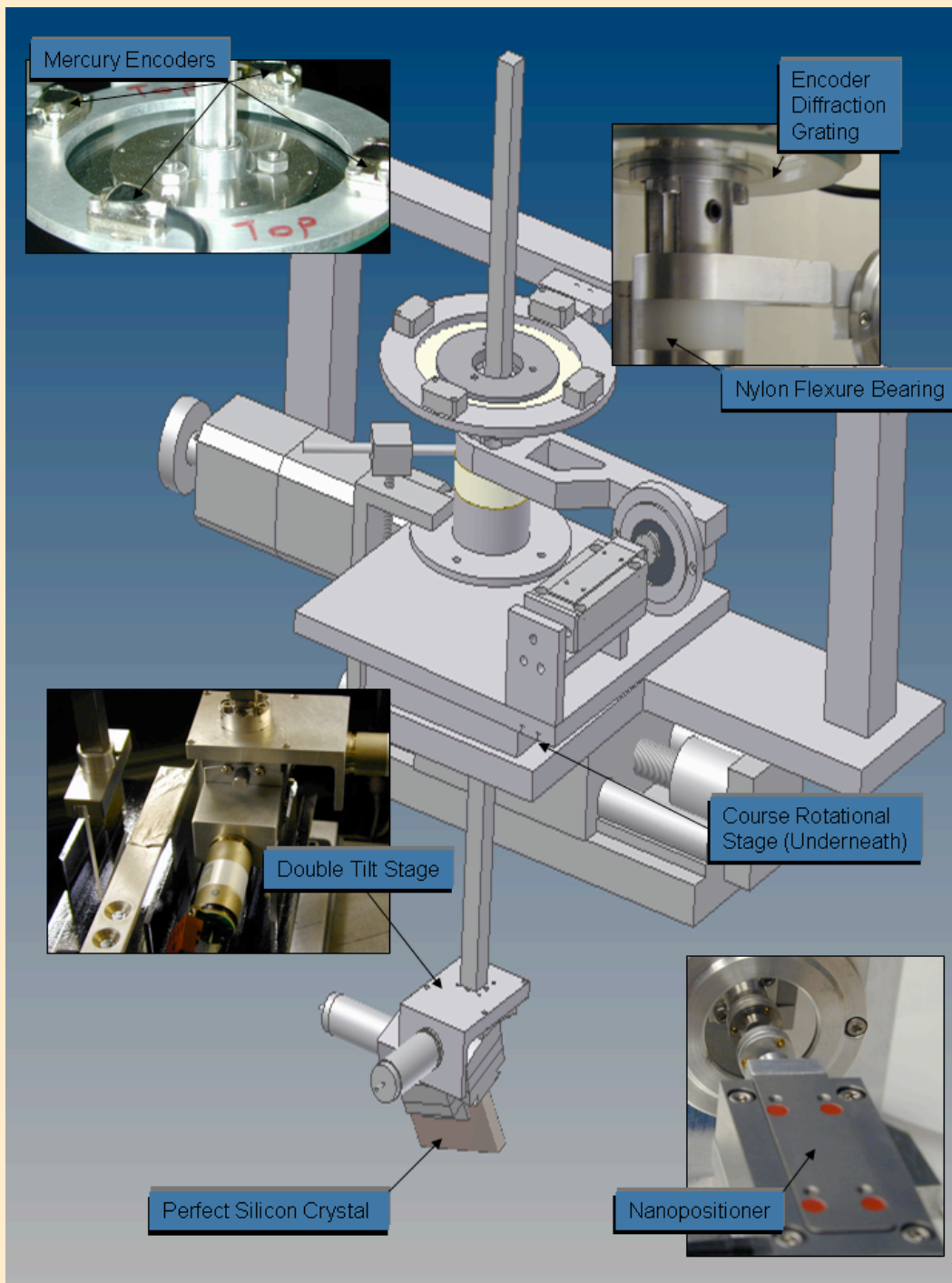
This will give b_{ne} , and hence $\langle r_n^2 \rangle$, to $< 1\%$

Tulane-NIST neutron charge radius experiment

10 cm lever with nylon
flexure bearing

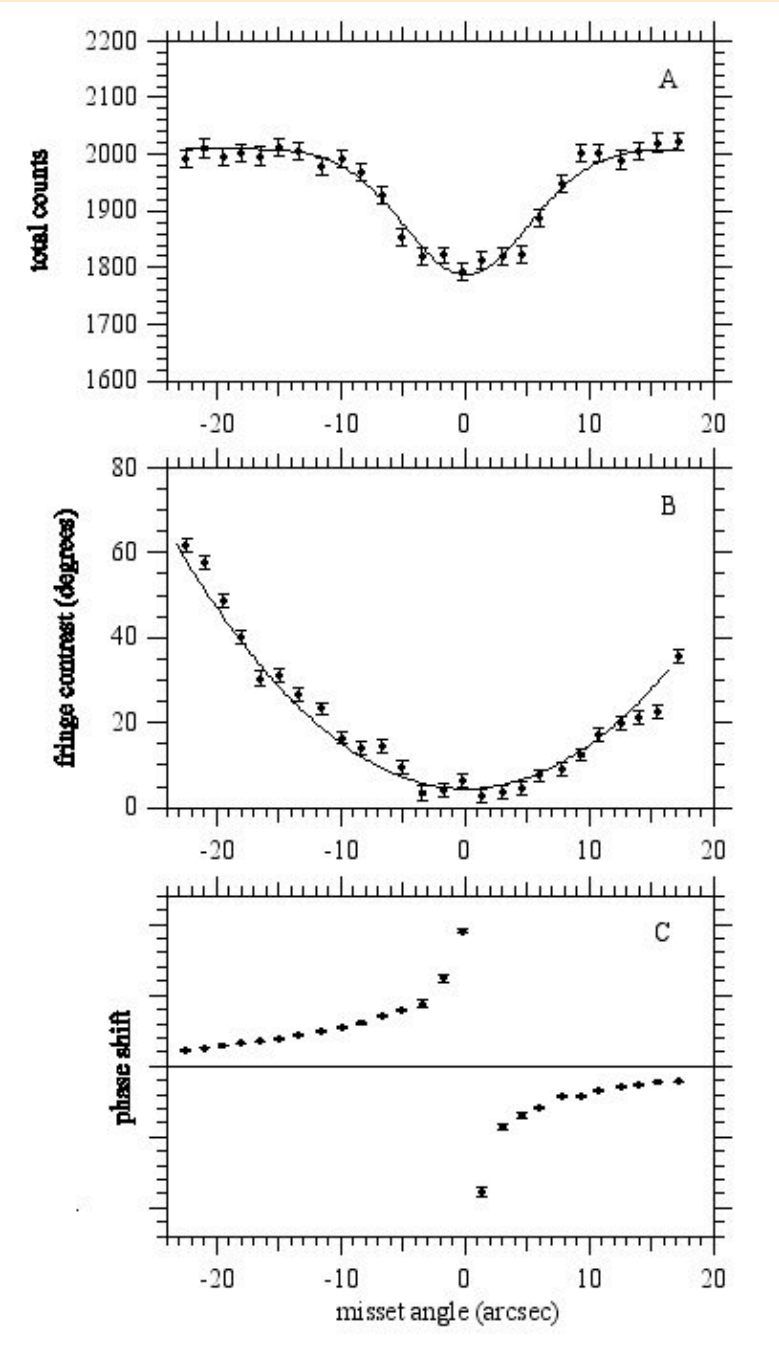
Physik Instrumente
P-753 PZT nanopositioner
25 μm range
1.0 nm precision (.002 arcsec)

Four Micro-E
mercury rotation encoders
.010 arcsec precision



Preliminary Data:

These data were taken at NIST in September 2005



The Neutron

by Gina Berkeley

When a pion an innocent proton seduces

With neither excuses

Abuses

Nor scorn

For its shameful condition

Without intermission

The proton produces: a neutron is born.

What love have you known

O neutron full grown

As you bombinate into the vacuum alone?

Its spin is $1/2$, and its mass is quite large

-about 1 AMU

but it hasn't a charge;

Though it finds satisfaction in strong interaction

It doesn't experience Coulombic attraction

But what can you borrow

Of love, joy, or sorrow

O neutron, when life has so short a tomorrow?

Within its

Twelve minutes

Comes disintegration

Which leaves an electron in mute desolation

And also another ingenuous proton

For other unscrupulous pions to dote on.

At last, a neutrino;

Alas, one can see no

Fulfilment for such a leptonic bambino.

No loving, no sinning

Just spinning and spinning

Eight times through the globe without ever beginning...

A cycle mechanic

No anguish or panic

For such is the pattern of life inorganic.

O better

The fret a

Poor human endures

Than the neutron's dichotic

Robotical

Amours.